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3.23 Electrical, Optical, and Magnetic Properties of Materials
Fall 2007

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3.23 Fall 2007 – Lecture 11

FERMI $\rightarrow \frac{e^2}{\hbar c}$

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Last time

1. Explicit solution of the Bloch equation, energy bands
2. Brillouin zone, Fermi surface
3. Energy of molecules and solids
4. Mean field approaches – Hartree and Hartree-Fock
5. Spin-statistics, Slater determinant, Pauli principle
6. Huckel approach (LCAO for aromatic compounds)

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Study

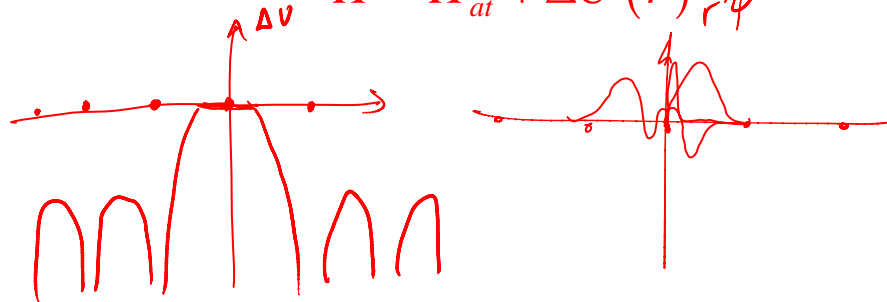
- Chap. 5 Singleton
- Read Chap. 6 Singleton

POST PROF. FINK → FREE ELECTRON GAS
FRIDAY 4pm 13-8066
MONDAY 4pm 13-8066

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Tight-binding (LCAO for solids)

- Hamiltonian $\hat{H} \stackrel{PER}{=} \hat{H}_{at} + \Delta \hat{U}(\vec{r}) \psi$



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Tight-binding (LCAO for solids)

- Bloch eigenstates of an ATOMIC CRYSTAL

$$\Psi_{n\vec{k}}(\vec{r}) = \sum_{\vec{R}} \exp(i\vec{k} \cdot \vec{R}) \psi_n(\vec{r} - \vec{R})$$

$$\begin{aligned} \Psi_{n\vec{k}}(\vec{r} + \vec{R}) &= \sum_{\vec{R}'} e^{i\vec{k} \cdot \vec{R}'} \psi_n(\vec{r} - \vec{R}' + \vec{R}) \\ &= e^{i\vec{k} \cdot \vec{R}} \sum_{\vec{R}'} e^{i\vec{k} \cdot (\vec{R}' - \vec{R})} \psi_n(\vec{r} - (\vec{R}' - \vec{R})) \end{aligned}$$

Bloch SUM

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Tight-binding (LCAO for solids)

- Bloch eigenstates of a REAL CRYSTAL

$$\Psi_{n\vec{k}}(\vec{r}) = \sum_{\vec{R}} \exp(i\vec{k} \cdot \vec{R}) \phi(\vec{r} - \vec{R})$$

$$\phi(\vec{r}) = \sum_n b_n \psi_n(\vec{r})$$

$$\int \psi_m^* (H_{AT} + UV) \psi_{nk} = E_n(k) \psi_{nk} \psi_m^*$$

LCAO

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Some despicable algebraic workout

$$\int \psi_m^* H_{\text{eff}} \psi_k = \int (H_{\text{eff}} \psi_m)^* \psi_k = E_m \int \psi_m^* \psi_k$$

$$\int \psi_m^* \Delta U \psi_k = (E(k) - E_m) \int \psi_m^* \psi_k$$

$$\begin{aligned} \int \psi_m^*(\vec{r}) \psi_k(\vec{r}) &= \sum_{\vec{R}} e^{i\vec{k} \cdot \vec{R}} \int \psi_m^*(\vec{r}) \left(\sum_n b_n \psi_n(\vec{r}-\vec{R}) \right) \\ &= b_m + \sum_{\vec{R} \neq 0} \sum_n b_n \int \psi_m^*(\vec{r}) \psi_n(\vec{r}-\vec{R}) e^{i\vec{k} \cdot \vec{R}} \end{aligned}$$

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More workout

$$\begin{aligned} (E(k) - E_m) b_m &= - (E(k) - E_m) \sum_n b_n \left(\int \psi_m^*(\vec{r}) \psi_n(\vec{r}-\vec{R}) e^{i\vec{k} \cdot \vec{R}} \right) \\ b_m \neq 0 \Rightarrow E(k) &\neq E_m \quad \left(\sum_{\vec{R} \neq 0} \int \psi_m^*(\vec{r}) \psi_n(\vec{r}-\vec{R}) e^{i\vec{k} \cdot \vec{R}} \right) \\ b_{n \neq m} &= 0 \quad + \sum_n b_n \int \psi_m^*(\vec{r}) \Delta U(\vec{r}) \psi_n(\vec{r}) \\ &\quad + \sum_n b_n \left(\sum_{\vec{R} \neq 0} \int \psi_m^*(\vec{r}) \Delta U(\vec{r}) \psi_n(\vec{r}-\vec{R}) \right) \end{aligned}$$

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More

$$\underbrace{\int \psi_m^*(r) \Delta U(r) \psi_m(r)}_{\text{by } = -\beta}$$

$$\underbrace{\int \psi_m^*(r) \psi_m(r - \vec{R})}_{\alpha(\vec{R})}$$

$$\underbrace{\int \psi_m^*(r) \Delta U(r) \psi_m(r - \vec{R})}_{-\mu(\vec{R})}$$

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More

$$E(k) - E_m = - \sum_{R \neq 0} e^{i\vec{k} \cdot \vec{R}} \alpha(\vec{R})$$

$$\Downarrow$$

$$E(k) = E_m - \beta \sum_{R \neq 0} \mu(\vec{R}) e^{i\vec{k} \cdot \vec{R}}$$
~~$$1 + \sum_{R \neq 0} \alpha(\vec{R}) e^{i\vec{k} \cdot \vec{R}}$$~~

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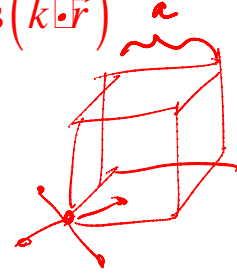
From s level to s bands

$$\varepsilon(\vec{k}) = E_s - \beta - \sum_{\text{nearest neighb.}} \gamma(\vec{R}) \cos(\vec{k} \cdot \vec{r})$$

$N a = \vec{R} = \frac{a}{2} (\pm 1, \pm 1, 0)$
 $(\pm 1, 0, \pm 1)$
 $(0, \pm 1, \pm 1)$

$\vec{k} = (k_x, k_y, k_z)$

$\varepsilon(\vec{k}) = E_s - \beta - 4\gamma \left(\cos \frac{1}{2} k_x a \cos \frac{1}{2} k_y a + \dots \right)$
ALL THIS PART!



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From s level to s bands

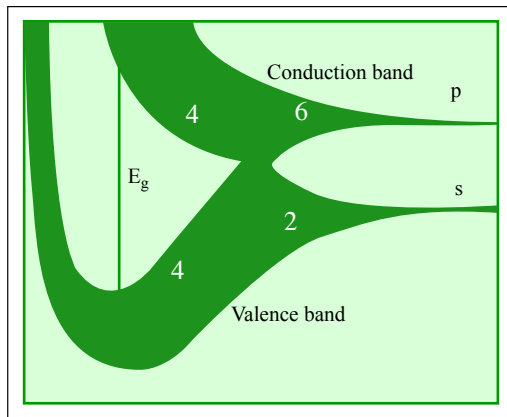


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Tight-binding vs. empirical psp

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Tight-binding vs. empirical psp

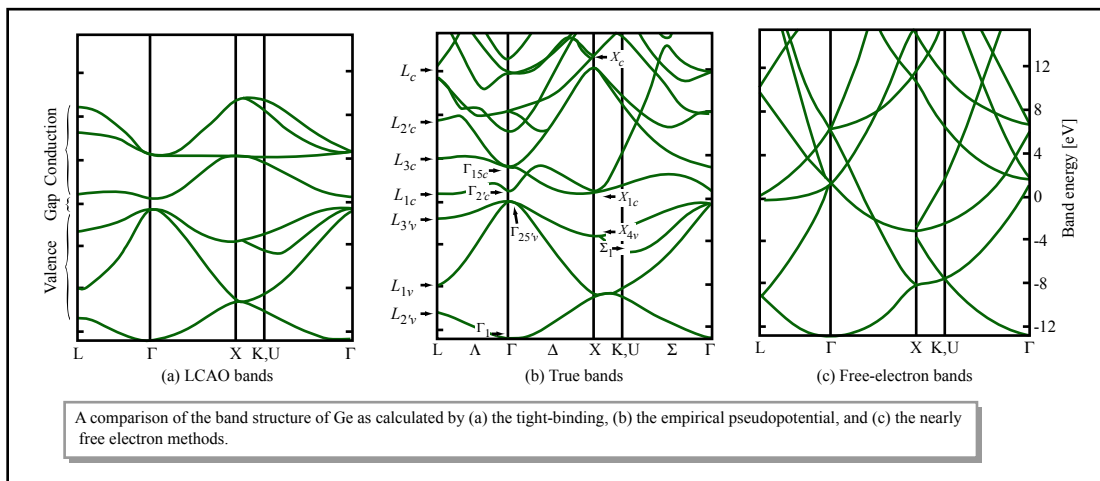


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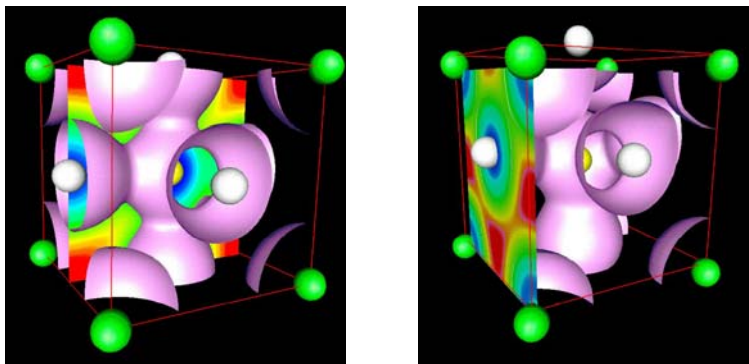
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Bands in Ge

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Ferroelectric perovskites



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Ferroelectric perovskites

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Ferroelectric perovskites

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