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3.23 Electrical, Optical, and Magnetic Properties of Materials
Fall 2007

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3.23 Fall 2007 – Lecture 2

THINK OUTSIDE THE BOX

More practical info

- Problem sets – out on Wed (and posted on Stellar), due by 5pm of the following weekend (after that 75%, after Thu 5pm 50%, after Fri 5pm 25%)
- ~11 in total, 30% of the grade
- Sometimes I mention homework – it's not the “Problem Set” @ Poilvert, Bonnet

Homework

- Take notes
- Revise posted lecture
- Study posted or assigned material
(TEXTBOOKS – do you have them ?)
- Meet with TAs or Instructor:

Last time: Wave mechanics

1. Particles, fields, and forces
2. Dynamics – from Newton to Schroedinger
3. De Broglie relation $\lambda \cdot p = h$
4. Waves and plane waves
5. Harmonic oscillator

Time-dependent Schrödinger's equation

(Newton's 2nd law for quantum objects)

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi(\vec{r}, t) + V(\vec{r}, t) \Psi(\vec{r}, t) = i\hbar \frac{\partial \Psi(\vec{r}, t)}{\partial t}$$

1925-onwards: E. Schrödinger (wave equation), W. Heisenberg (matrix formulation), P.A.M. Dirac (relativistic)

Plane waves as free particles

Our free particle $\Psi(\vec{r}, t) = A \exp[i(\vec{k} \cdot \vec{r} - \omega t)]$ satisfies the wave equation:

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi(\vec{r}, t) = i\hbar \frac{\partial \Psi(\vec{r}, t)}{\partial t} \quad \left(\text{provided } E = \hbar\omega = \frac{p^2}{2m} = \frac{\hbar^2 k^2}{2m} \right)$$

Stationary Schrödinger's Equation (I)

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi(\vec{r}, t) + V(\vec{r}, *) \Psi(\vec{r}, t) = i\hbar \frac{\partial \Psi(\vec{r}, t)}{\partial t}$$

Stationary Schrödinger's Equation (II)

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \varphi(\vec{r}) = E \varphi(\vec{r})$$

Stationary Schrödinger's Equation (III)

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \varphi(\vec{r}) = E \varphi(\vec{r})$$

1. It's not proven – it's postulated, and it is confirmed experimentally
2. It's an “eigenvalue” equation: it has a solution only for certain values (discrete, or continuum intervals) of E
3. For those eigenvalues, the solution (“eigenstate”, or “eigenfunction”) is the complete descriptor of the electron in its equilibrium ground state, in a potential $V(r)$.
4. As with all differential equations, boundary conditions must be specified
5. Square modulus of the wavefunction = probability of finding an electron

Free particle: $\Psi(x,t)=\varphi(x)f(t)$

$$-\frac{\hbar^2}{2m}\nabla^2\varphi(x) = E\varphi(x)$$



$$i\hbar\frac{d}{dt}f(t) = Ef(t)$$



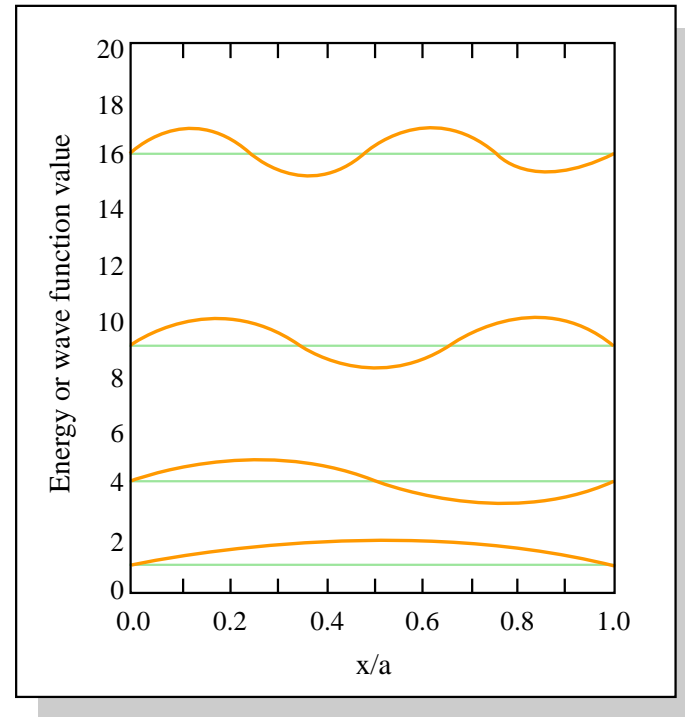
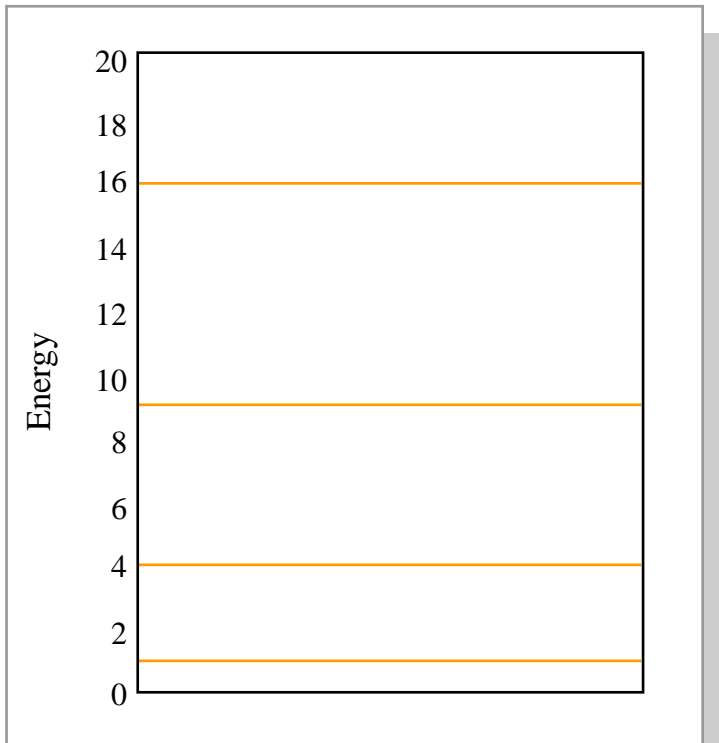
Infinite Square Well (I)

(particle in a 1-dim box)

$$-\frac{\hbar^2}{2m} \frac{d^2 \varphi(x)}{dx^2} = E \varphi(x)$$

Infinite Square Well (II)

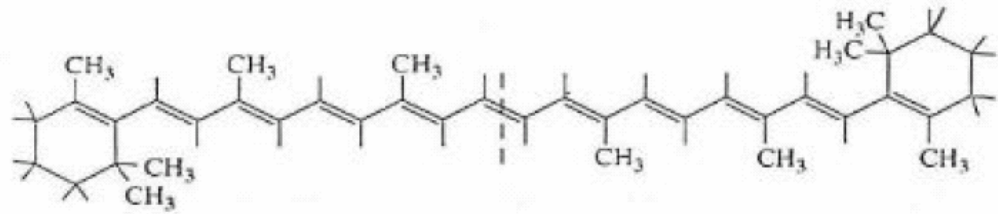
Infinite Square Well (III)



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The power of carrots

- β -carotene



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Physical Observables from Wavefunctions

- Eigenvalue equation:

$$\left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \varphi(x) = E\varphi(x)$$

- Expectation values for the operator (energy)

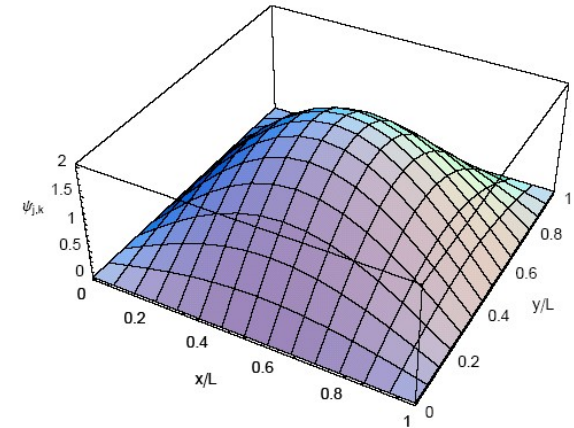
$$E = \int \varphi^*(x) \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \varphi(x) dx$$

Particle in a 2-dim box

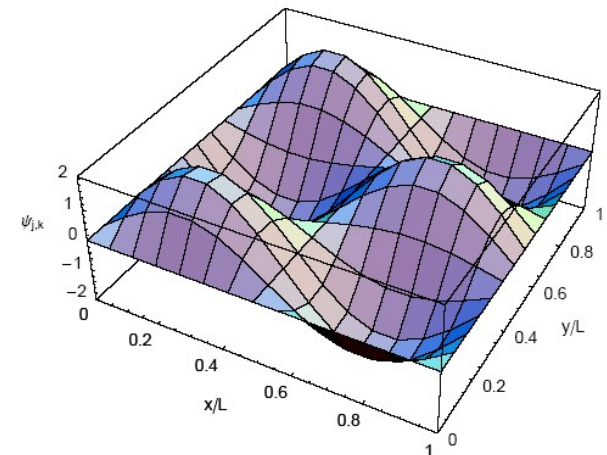
$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \varphi(x, y) = E \varphi(x, y)$$

Particle in a 2-dim box

$$\varphi(x, y) = C \sin\left(\frac{l\pi x}{a}\right) \sin\left(\frac{m\pi y}{b}\right)$$



$$E = \frac{h^2}{8m} \left(\frac{l^2}{a^2} + \frac{m^2}{b^2} \right)$$



Particle in a 3-dim box

$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \varphi(x, y, z) = E \varphi(x, y, z)$$

Particle in a 3-dim box: *Farbe* defect in halides (e⁻ bound to a negative ion vacancy)

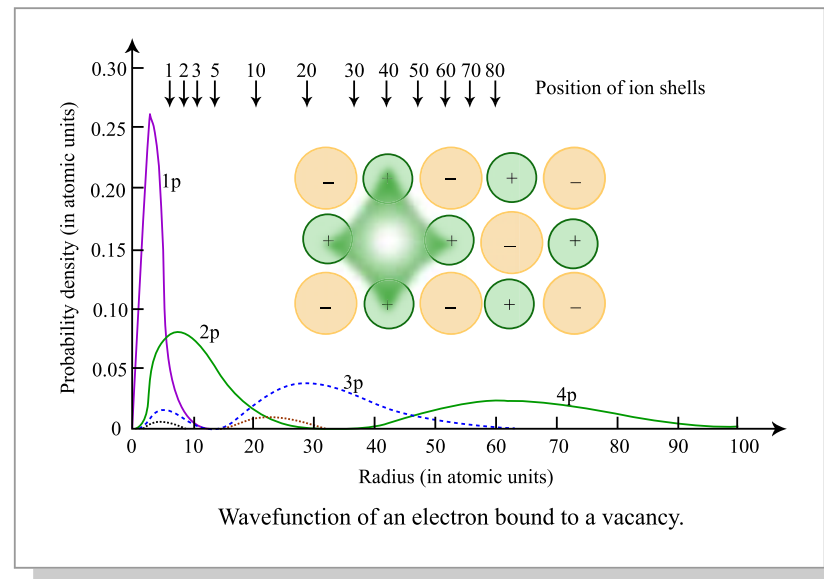
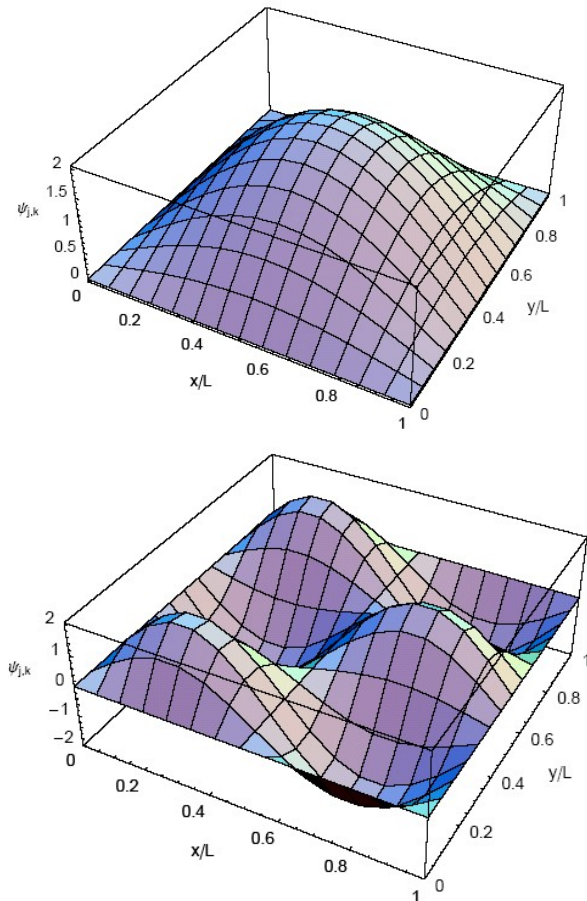


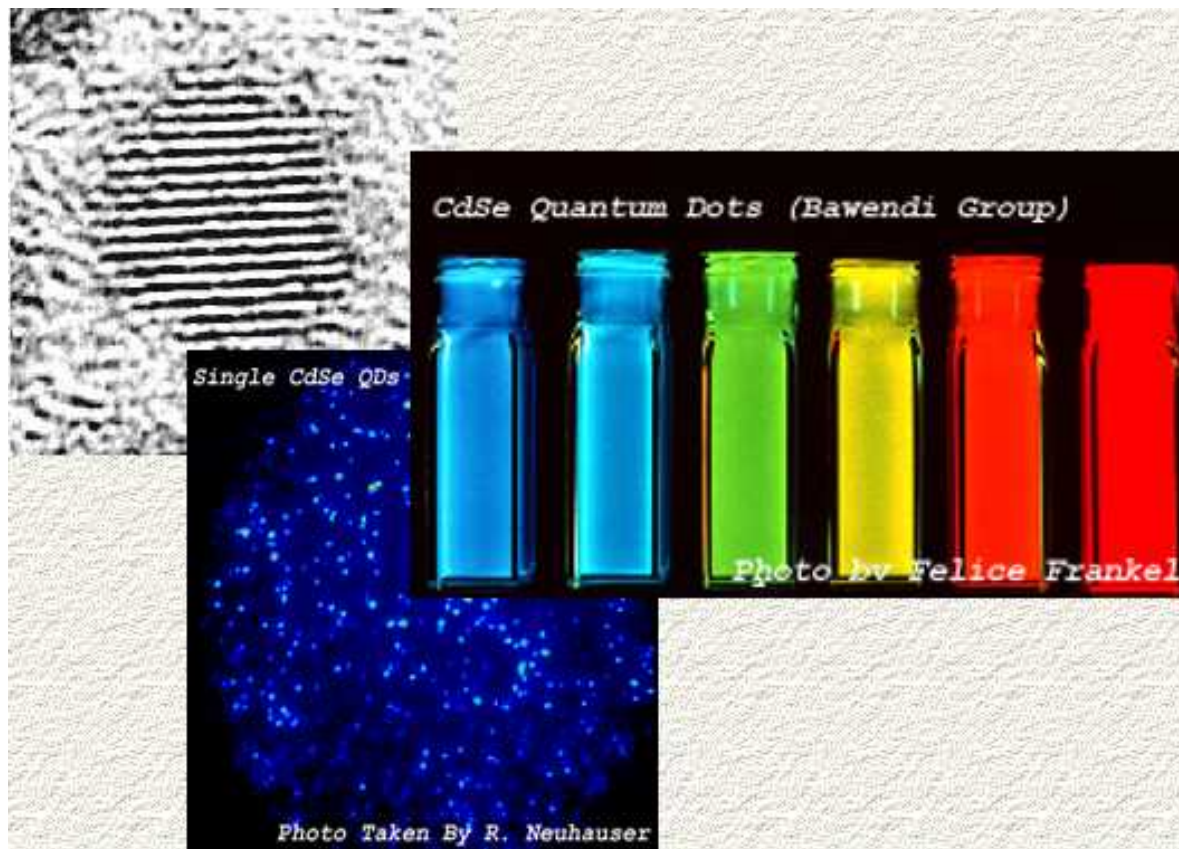
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From Carl Zeiss to MIT...

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Please see Avakian, P., and Smakula, A. "Color Centers in Cesium Halide Single Crystals." *Physical Review* 120 (December 1960): 2007.

Light absorption/emission



Courtesy M. Bawendi and Felice Frankel.
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Abstract and Fig. 1 in Willey, T. M., et al. "Molecular Limits to the Quantum
Confinement Model in Diamond Clusters." *Physical Review Letters* 95 (2005): 113401.

Metal Surfaces (I)

$$\left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \varphi(x) = E\varphi(x)$$

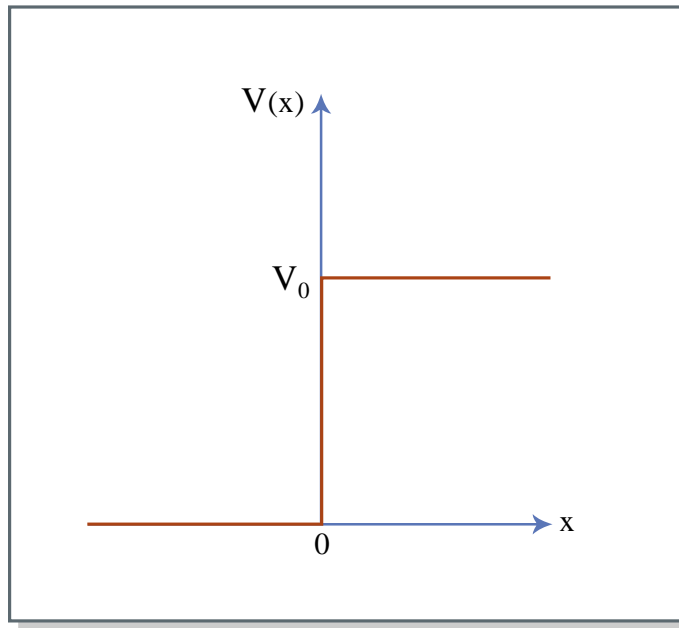


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Metal Surfaces (II)

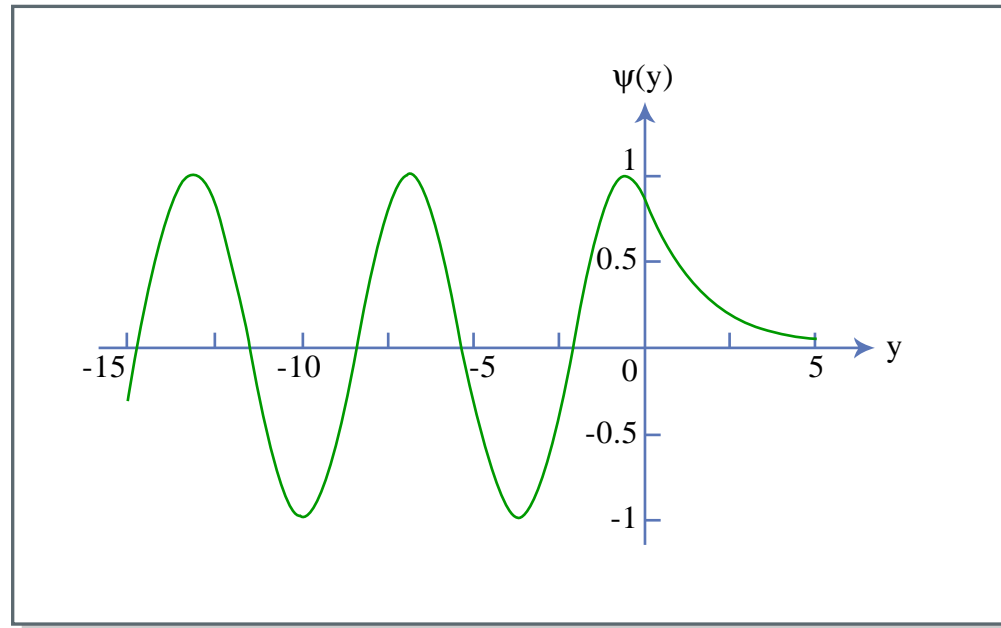
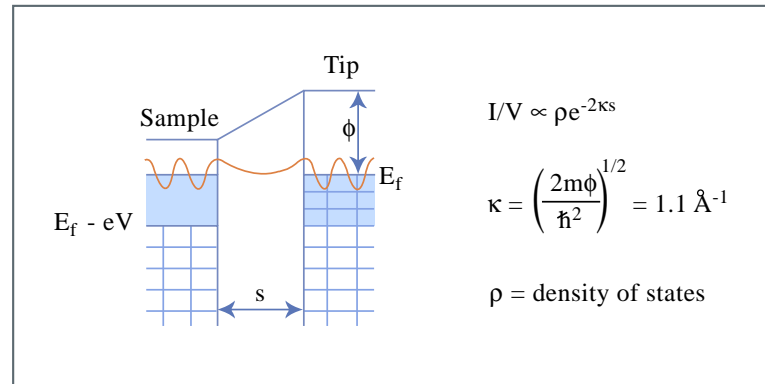
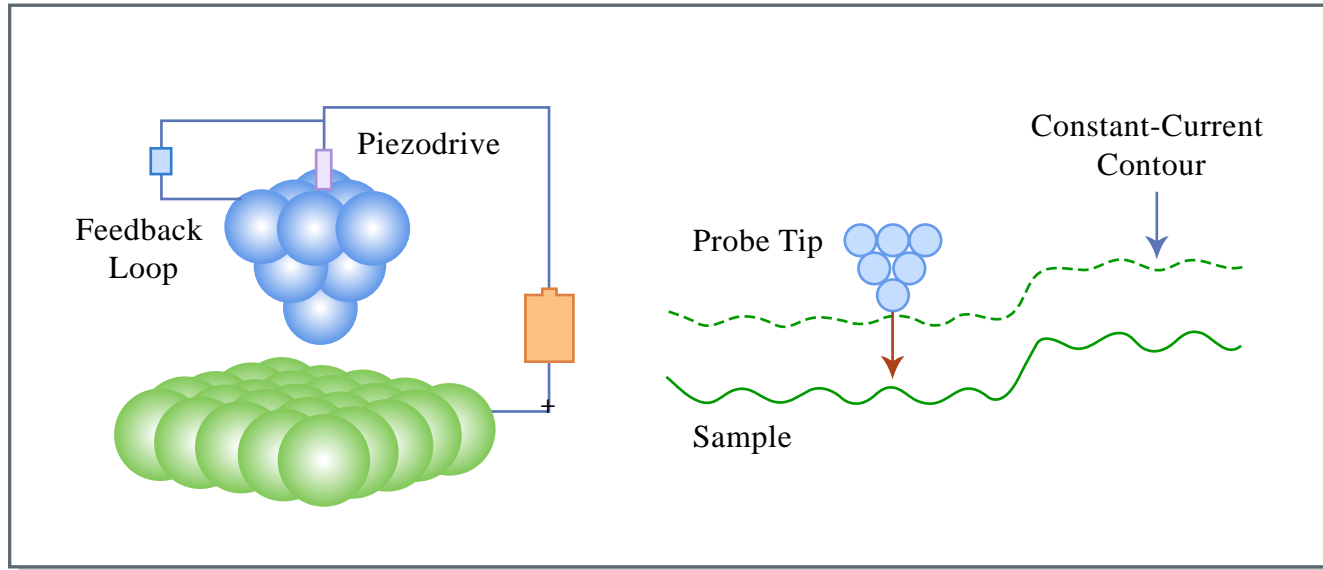


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Scanning Tunnelling Microscopy



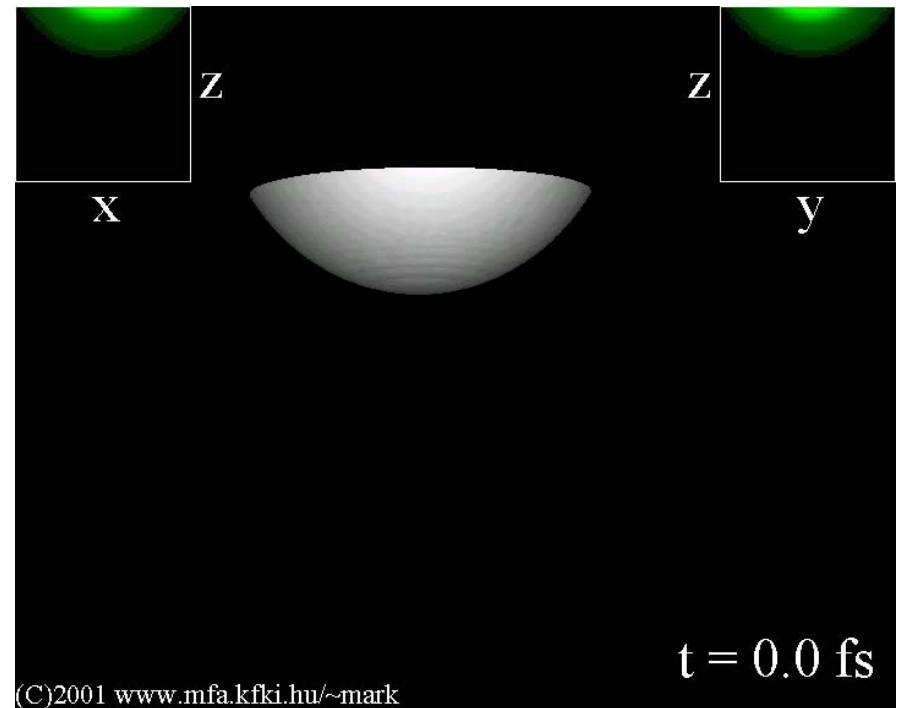
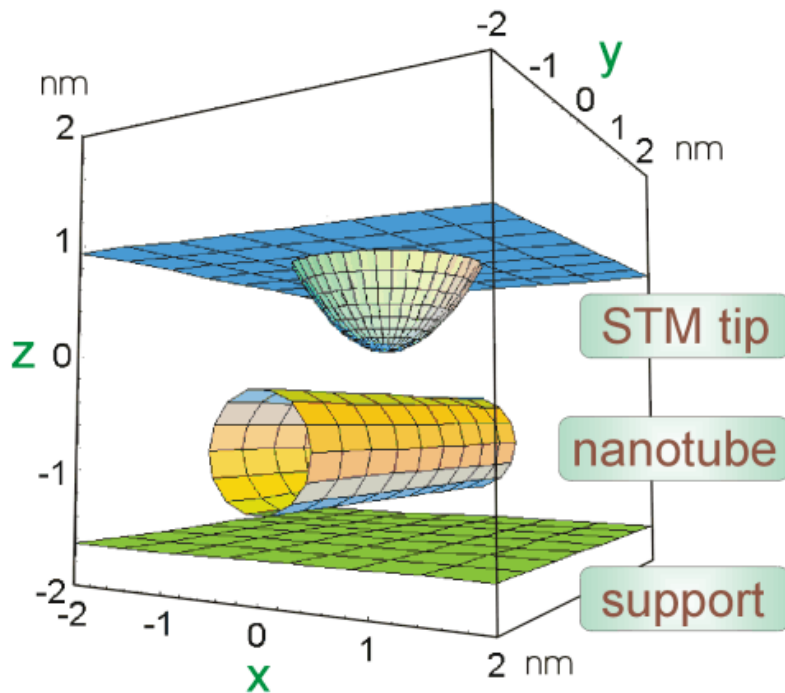
$$I/V \propto \rho e^{-2\kappa s}$$

$$\kappa = \left(\frac{2m\phi}{\hbar^2} \right)^{1/2} = 1.1 \text{ \AA}^{-1}$$

ρ = density of states

Figures by MIT OpenCourseWare.

Wavepacket tunnelling through a nanotube



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<http://newton.phy.bme.hu/education/schrd/index.html>

<http://www.quantum-physics.polytechnique.fr>