

Honeycombs - In-plane behaviour

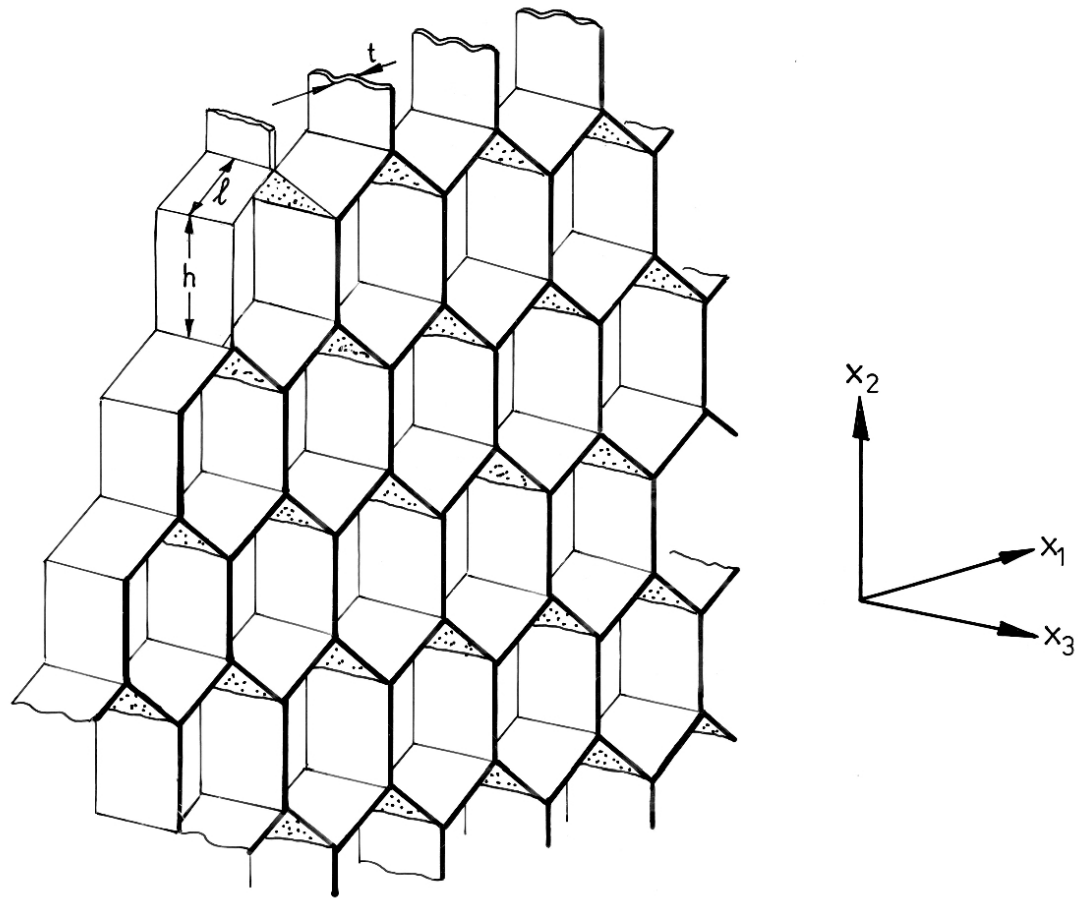
- prismatic cells
 - polymer, metal, ceramic honeycombs widely available
 - used for sandwich structure cores, energy absorption, carriers for catalysts
 - some natural materials (eg. wood, cork) can be idealized as honeycombs
 - mechanisms of deformation + failure in hexagonal honeycombs parallel those in foams
 - simpler geometry (unit cell) - easier to analyze
 - mechanisms of deformation in triangular honeycombs parallels those in 3D trusses (lattice materials)
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Stress - Strain curves + Deformation behaviour : In - Plane

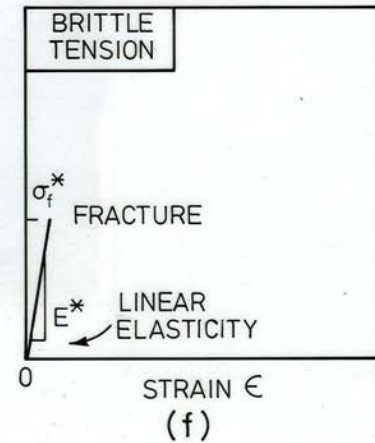
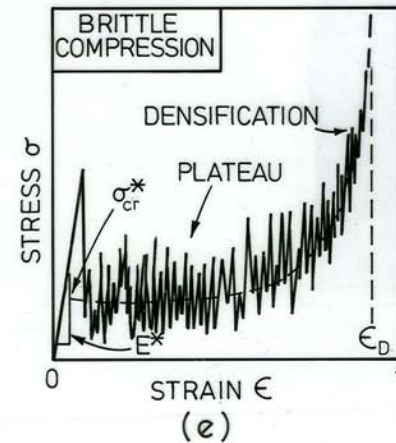
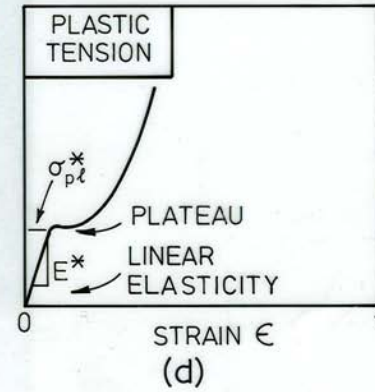
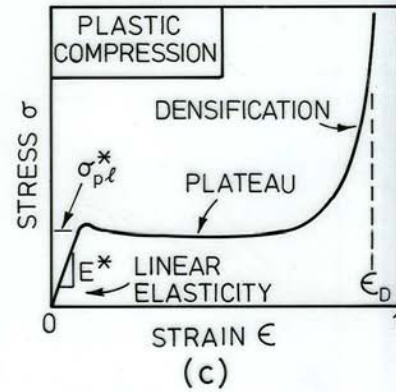
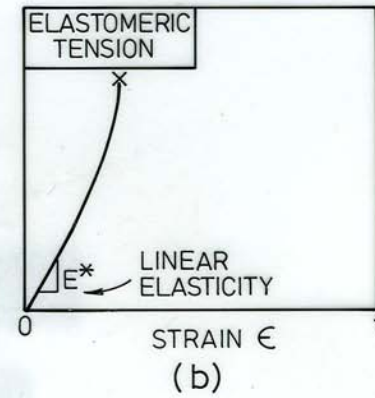
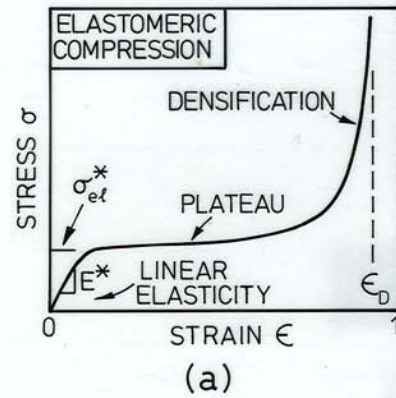
Compression

- 3 regimes - linear elastic - bending
 - stress plateau - buckling
 - yielding
 - brittle crushing
- densification - cell walls touch
- increasing $t/l \Rightarrow E^* \uparrow \sigma^* \uparrow \epsilon_D \downarrow$

Honeycomb Geometry

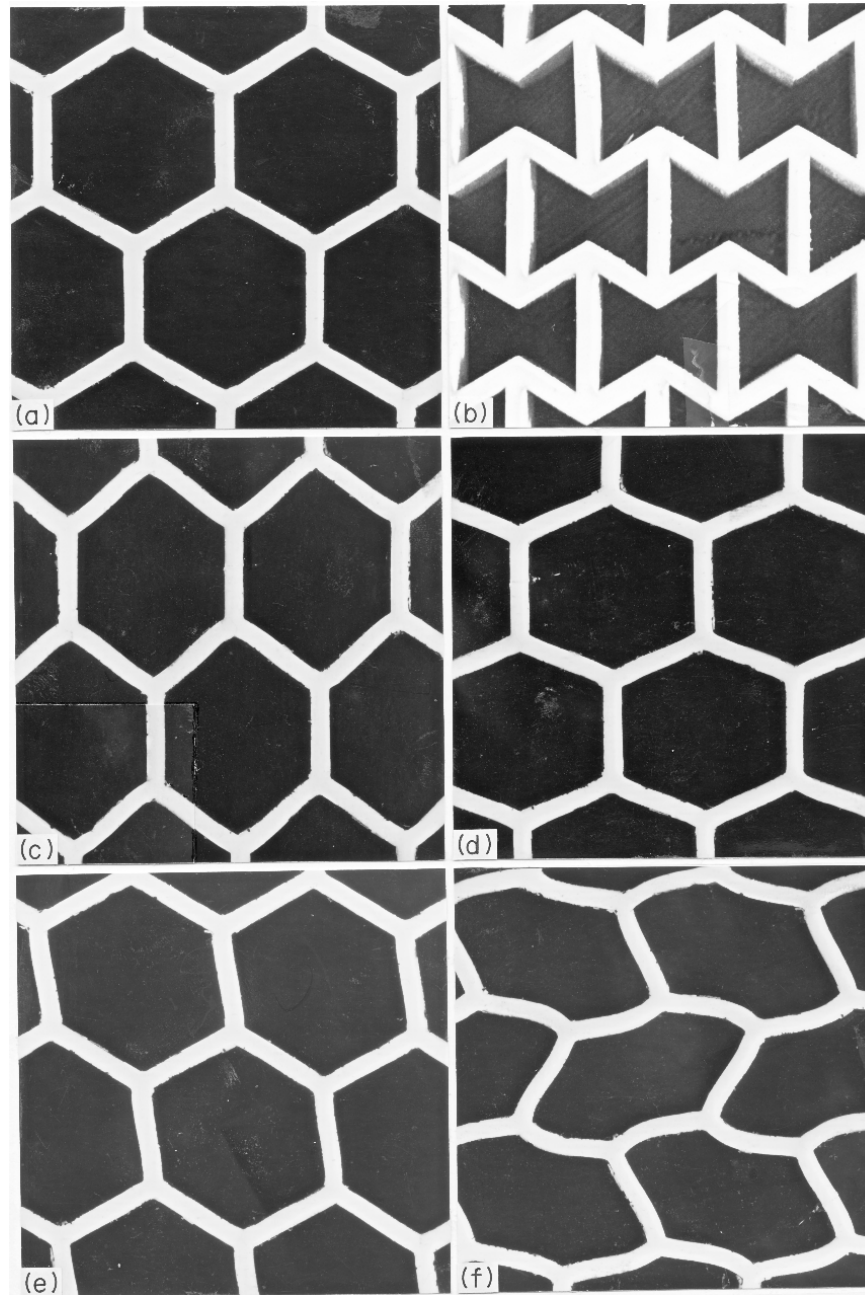


Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.



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Deformation mechanisms

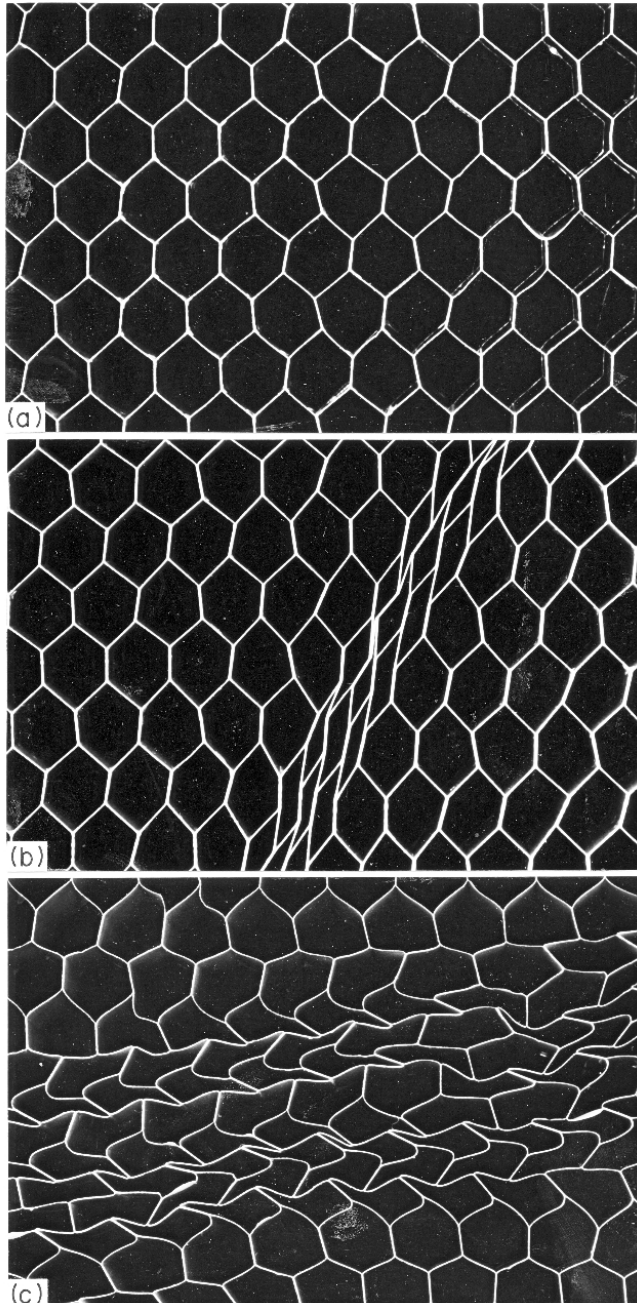


Bending
 X_1 Loading

Bending
 X_2 Loading

Bending
Shear

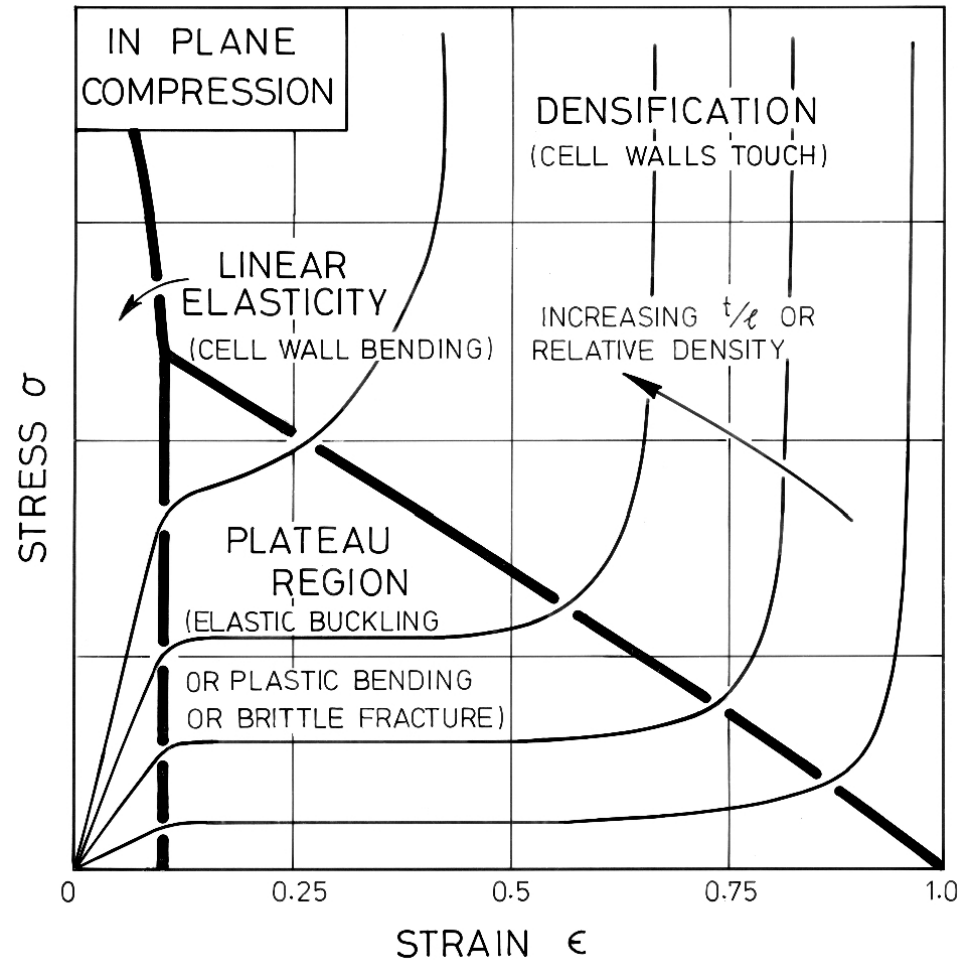
Buckling



Plastic collapse in an aluminum honeycomb

Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

Stress-Strain Curve



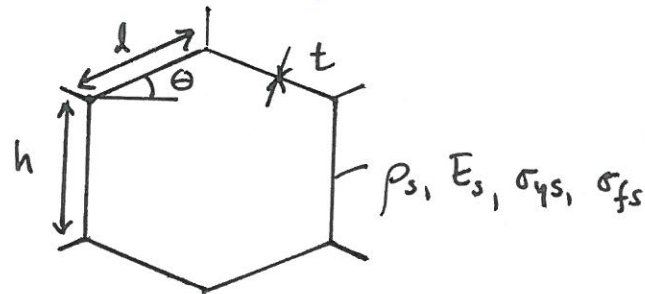
Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

tension

- linear elastic - bending
- stress plateau - exists only if cell walls yield
 - no buckling in tension
 - brittle honeycombs fracture in tension

Variables affecting honeycomb properties

- relative density $\frac{\rho^*}{\rho_s} = \frac{t/l (h/l + 2)}{2 \cos \theta (h/l + \sin \theta)} = \frac{2}{\sqrt{3}} \frac{t}{l}$ regular hexagons
- solid cell wall properties: $\rho_s, E_s, \sigma_{ys}, \sigma_{fs}$
- cell geometry: $h/l, \theta$



In-plane properties

assumptions:

- ρ^*/ρ_s small - neglect axial + shear contribution to def^{ion}
- deformations small - neglect changes in geometry
- cell wall - linear elastic, isotropic

symmetry

- honeycombs are orthotropic - rotate 180° about each of three mutually perpendicular axes & structure is the same.

Linear elastic deformation

$$\begin{bmatrix} \epsilon_1 \\ \epsilon_2 \\ \epsilon_3 \\ \epsilon_4 \\ \epsilon_5 \\ \epsilon_6 \end{bmatrix} = \begin{bmatrix} 1/E_1 & -\nu_{21}/E_2 & -\nu_{31}/E_3 & 0 & 0 & 0 \\ -\nu_{12}/E_1 & 1/E_2 & -\nu_{32}/E_3 & 0 & 0 & 0 \\ -\nu_{13}/E_1 & -\nu_{23}/E_2 & 1/E_3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1/G_{23} & 0 & 0 \\ 0 & 0 & 0 & 0 & 1/G_{13} & 0 \\ 0 & 0 & 0 & 0 & 0 & 1/G_{12} \end{bmatrix} \begin{bmatrix} \sigma_1 \\ \sigma_2 \\ \sigma_3 \\ \sigma_4 \\ \sigma_5 \\ \sigma_6 \end{bmatrix}$$

(Symmetric)

- Matrix notation:

$\epsilon_1 = \epsilon_{11}$	$\epsilon_4 = \gamma_{23}$	$\sigma_1 = \sigma_{11}$	$\sigma_4 = \sigma_{23}$
$\epsilon_2 = \epsilon_{22}$	$\epsilon_5 = \gamma_{13}$	$\sigma_2 = \sigma_{22}$	$\sigma_5 = \sigma_{13}$
$\epsilon_3 = \epsilon_{33}$	$\epsilon_6 = \gamma_{12}$	$\sigma_3 = \sigma_{33}$	$\sigma_6 = \sigma_{12}$

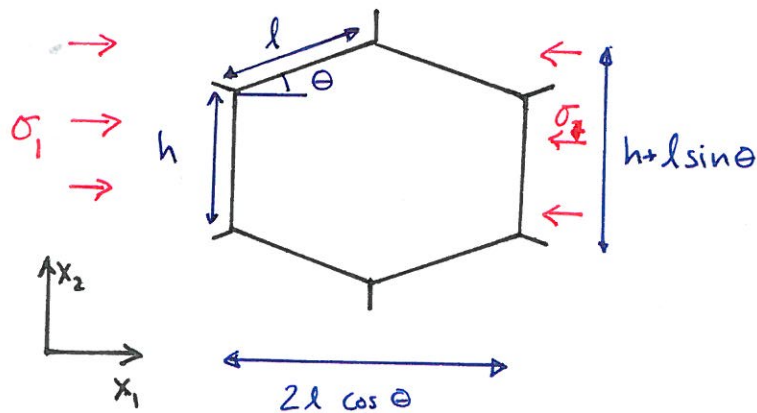
- in-plane (x_1-x_2): 4 independent elastic constants:

$$E_1, E_2, \nu_{12}, G_{12}$$

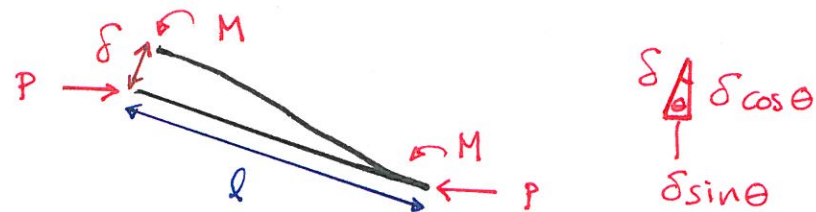
and compliance matrix symmetric - $\frac{\nu_{12}}{E_1} = -\frac{\nu_{21}}{E_2}$ (reciprocal relation)

[notation for Poisson's ratio: $\nu_{ij} = -\frac{\epsilon_j}{\epsilon_i}$]

Young's modulus in x_1 direction



unit cell in x_1 direction: $2l \cos \theta$
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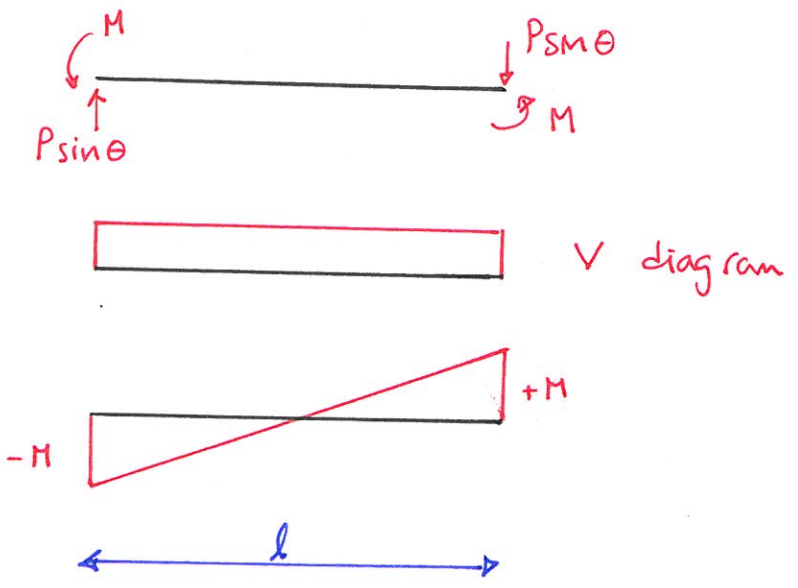


$$\sigma_1 = \frac{P}{(h + l \sin \theta) b}$$

$$\epsilon_1 = \frac{\delta \sin \theta}{l \cos \theta}$$

In-Plane Deformation: Linear Elasticity

Figure removed due to copyright restrictions. See Figure 5: L. J. Gibson,
M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials.](#)"



M diagram: 2 cantilevers of length $l/2$

$$\delta = 2 \cdot \frac{P \sin \theta (l/2)^3}{3 E_s I}$$

$$= \frac{2 P l^3 \sin \theta}{24 E_s I}$$

$$\delta = \frac{P l^3 \sin \theta}{12 E_s I}$$

$$I = \frac{b t^3}{12}$$

Combining: $E_1^* = \frac{\sigma_1}{\epsilon_1} = \frac{P}{(h + l \sin \theta) b} \frac{l \cos \theta}{\delta \sin \theta}$

$$= \frac{P}{(h + l \sin \theta) b} \frac{l \cos \theta}{l^3 \sin^2 \theta} \cancel{1/2} E_s \cancel{1/2} \frac{t^3}{1/2}$$

$$E_1^* = E_s \left(\frac{t}{l}\right)^3 \frac{\cos \theta}{(h/l + \sin \theta) \sin^2 \theta}$$

solid property relative density cell geometry

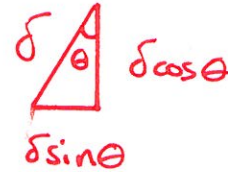
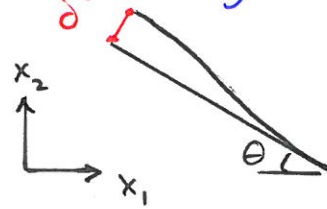
$$= \frac{4}{\sqrt{3}} \left(\frac{t}{l}\right)^3 E_s \quad \text{regular hexagons}$$

$h/l = 1 \quad \theta = 30^\circ$

6

Poisson's ratio for loading in x_1 direction

$$\nu_{12}^* = - \frac{\epsilon_2}{\epsilon_1}$$



$$\epsilon_1 = - \frac{\delta \sin \theta}{l \cos \theta} \text{ (shortens)}$$

$$\epsilon_2 = \frac{\delta \cos \theta}{h + l \sin \theta} \text{ (lengthens)}$$

$$\nu_{12}^* = - \frac{\delta \cos \theta}{h + l \sin \theta} \left(- \frac{l \cos \theta}{\delta \sin \theta} \right) = \frac{\cos^2 \theta}{(h/l + \sin \theta) \sin \theta}$$

- ν_{12}^* depends only on cell geometry ($h/l, \theta$), not on $E_s, t/l$
- regular hexagonal cells: $\nu_{12}^* = 1$
- ν can be negative for $\theta < 0$

eg. $h/l = 2 \quad \theta = -30^\circ \quad \nu_{12}^* = \frac{3/4}{(3/2)(-1/2)} = -1$

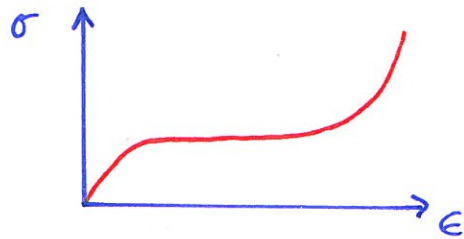
$E_2^* \quad \nu_{21}^* \quad G_{12}^*$

- can be found in similar way; results in book.

Compressive strength (plateau stress)

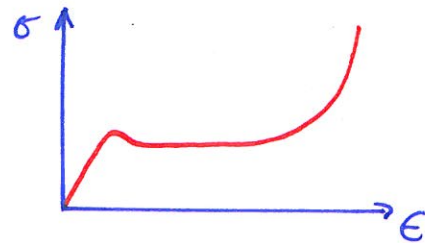
• cell collapse by:

(1) elastic buckling



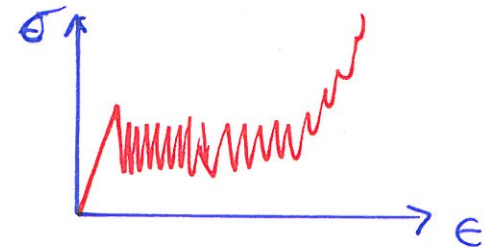
- buckling of vertical struts throughout honeycomb

(2) plastic yielding



- localization of yield
- as def^{ion} progresses, propagation of failure band

(3) brittle crushing

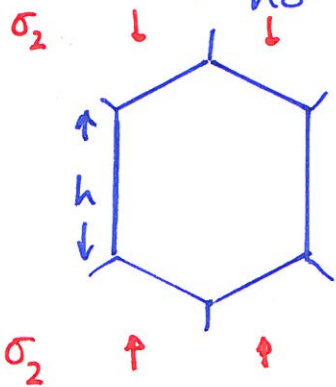


- peaks + valleys correspond to fracture of individual cell walls

Plateau stress : elastic buckling, σ_{el}^*

- elastomeric honeycombs - cell collapse by elastic buckling of walls of length h when loaded in x_2 direction

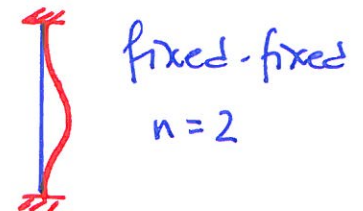
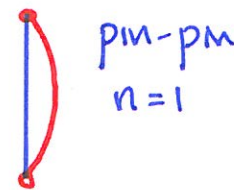
- no buckling for σ_1 ; bending of inclined walls goes to densification.



Euler buckling load

$$P_{cr} = \frac{n^2 \pi^2 E_s I}{h^2}$$

n = end constraint factor



Elastic Buckling

Figure removed due to copyright restrictions. See Figure 7: L. J. Gibson, M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials](#)."

- here, constraint n depends on stiffness of adjacent inclined members
- can find by elastic line analysis (see appendix if interested)
- rotational stiffness at ends of column, h , matched to rotational stiffness of inclined members
- find $h/l = 1 \quad 1.5 \quad 2$
 $n = 0.686 \quad 0.760 \quad 0.860$

and $(\sigma_{el}^*)_2 = \frac{P_{cr}}{2l \cos \theta b} = \frac{n^2 \pi^2 E_s}{h^2 2l \cos \theta b} \frac{bt^3}{12}$

$$(\sigma_{el}^*)_2 = \frac{n^2 \pi^2}{24} E_s \frac{(t/l)^3}{(h/l)^2 \cos \theta}$$

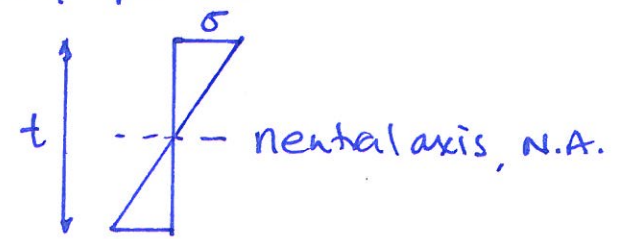
regular hexagons: $(\sigma_{el}^*)_2 = 0.22 E_s (t/l)^3$

‡ since $E_2^* = 4/\sqrt{3} E_s (t/l)^3 = 231 E_s (t/l)^3$

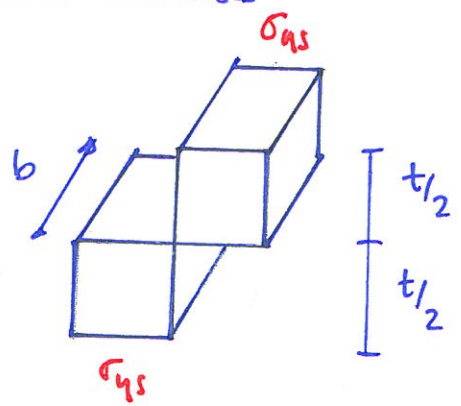
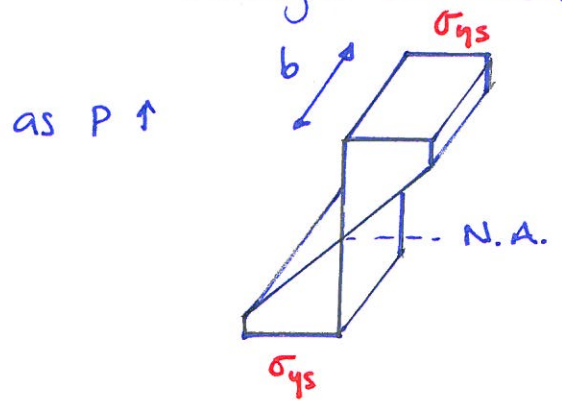
strain at buckling $(\epsilon_{el}^*)_2 = 0.10$, for regular hexagons independent of $E_s, t/l$

Plateau stress: plastic yielding, σ_{pl}^*

- failure by yielding in cell walls
- yield strength of cell walls = σ_{ys}
- plastic hinge forms when cross-section fully yielded
- beam theory - linear elastic $\sigma = \frac{My}{I}$



- once stress at outer fiber = σ_{ys} , yielding begins & then progresses through the section, as the load increases



- When section fully yielded (right fig.), form plastic "hinge"
- section rotates, like a pin

Plastic Collapse

Figure removed due to copyright restrictions. See Figure 8: L. J. Gibson, M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials](#)."

- moment at formation of plastic hinge (plastic moment, M_p):

$$M_p = \left(\sigma_{ys} b \frac{t}{2} \right) \left(\frac{t}{2} \right) = \frac{\sigma_{ys} b t^2}{4}$$

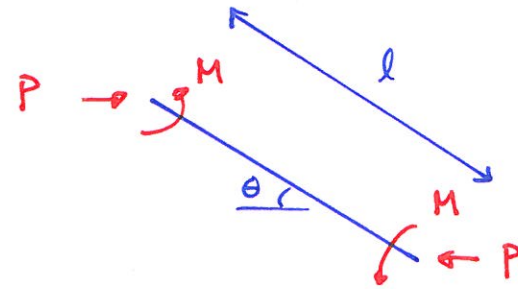
- applied moment, from applied stress

$$2 M_{app} - P l \sin \theta = 0$$

$$M_{app} = \frac{P l \sin \theta}{2}$$

$$\sigma_1 = \frac{P}{(h + l \sin \theta) b}$$

$$M_{app} = \sigma_1 (h + l \sin \theta) b \frac{l \sin \theta}{2}$$



- plastic collapse of honeycomb at $(\sigma_{pl}^*)_1$, when $M_{app} = M_p$

$$(\sigma_{pl}^*)_1 (h + l \sin \theta) \frac{l \sin \theta}{2} = \sigma_{ys} \frac{b t^2}{4}$$

$$\boxed{(\sigma_{pl}^*)_1 = \sigma_{ys} \left(\frac{t}{l} \right)^2 \frac{1}{2 (h/l + \sin \theta) \sin \theta}}$$

regular hexagons: $(\sigma_{pl}^*)_1 = \frac{2}{3} \sigma_{ys} \left(\frac{t}{l} \right)^2$

similarly, $(\sigma_{pl}^*)_2 = \sigma_{ys} \left(\frac{t}{l} \right)^2 \frac{1}{2 \cos^2 \theta}$

• can do similar analysis for ...

- for thin-walled honeycombs, elastic buckling can precede plastic collapse (for σ_2)
- elastic buckling stress = plastic collapse stress $(\sigma_{el}^*)_2 = (\sigma_{pl}^*)_2$

$$\frac{n^2 \pi^2 E_s}{24} \frac{(t/l)^3}{(h/l)^2 \cos \theta} = \frac{\sigma_{ys} (t/l)^2}{2 \cos^2 \theta}$$

$$(t/l)_{\text{critical}} = \frac{12 (h/l)^2}{n^2 \pi^2 \cos \theta} \left(\frac{\sigma_{ys}}{E_s} \right)$$

regular hexagons: $(t/l)_{\text{critical}} = 3 \frac{\sigma_{ys}}{E_s}$

- e.g. metals $\sigma_{ys}/E_s \sim .002$ $(t/l)_{\text{crit}} \sim 0.6\%$
 - most metal honeycomb denser than this
- polymers $\sigma_{ys}/E_s \sim 3-5\%$ $(t/l)_{\text{crit}} \sim 10-15\%$
 - low density polymers with yield point may buckle before yield.

Plateau stress: brittle crushing, (σ_c^*) ,

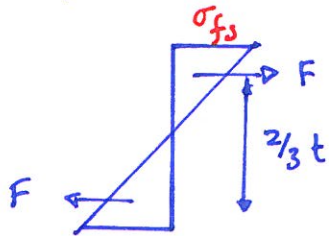
- ceramic honeycombs - fail in brittle manner
- cell wall bending - stress reaches modulus of rupture - wall fracture

loading in x_1 direction:

$$P = \sigma_1 (h + l \sin \theta) b \quad \sigma_{fs} = \text{modulus of rupture of cell wall}$$

$$M_{\text{max applied}} = \frac{Pl \sin \theta}{2} = \frac{\sigma_1 (h + l \sin \theta) b l \sin \theta}{2}$$

Moment at fracture, M_f

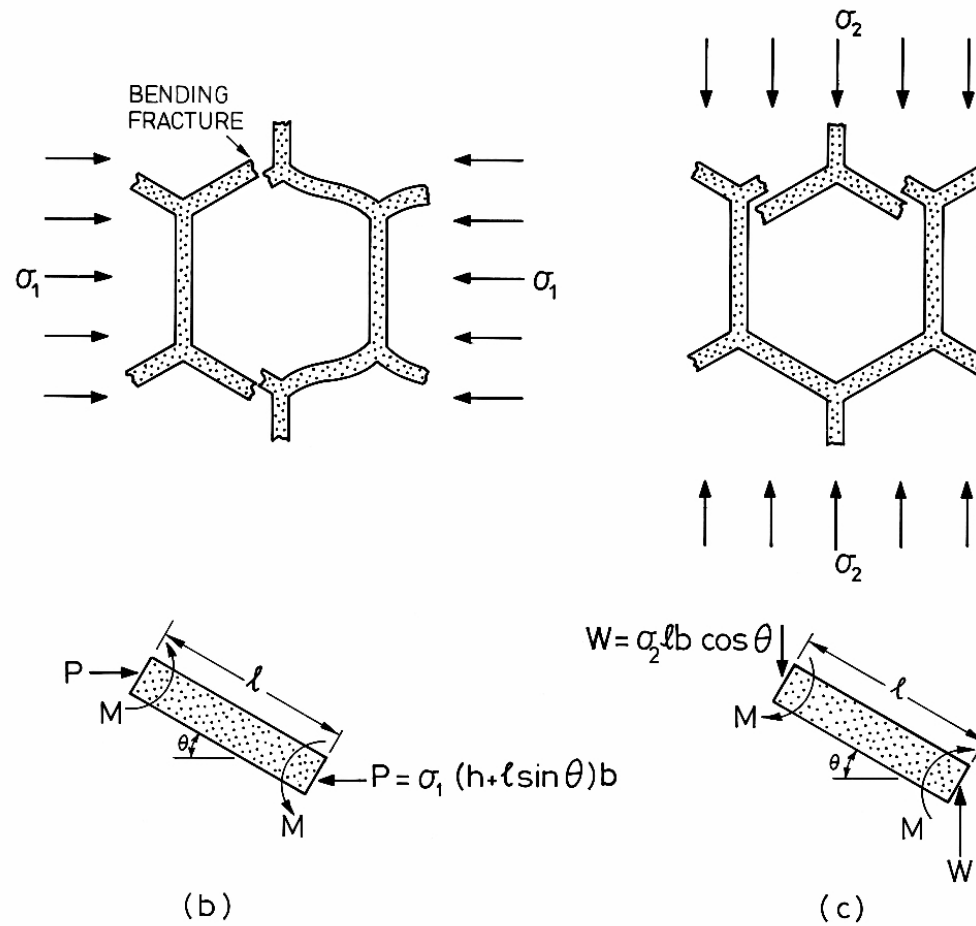


$$M_f = \left(\frac{1}{2} \sigma_{fs} b \frac{t}{2} \right) \left(\frac{2}{3} t \right) = \frac{\sigma_{fs} b t^2}{6}$$

$$\boxed{(\sigma_c^*)_1 = \sigma_{fs} \left(\frac{t}{l} \right)^2 \frac{1}{3 (h/l + \sin \theta) \sin \theta}}$$

regular hexagons: $(\sigma_c^*)_1 = \frac{4}{9} \sigma_{fs} \left(\frac{t}{l} \right)^2$

Brittle Crushing



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

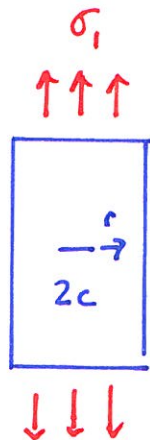
Tension

- no elastic buckling
- plastic plateau stress approx same in tension + compression (small geometric difference due to deformation)
- brittle honeycombs: fast fracture

Fracture toughness

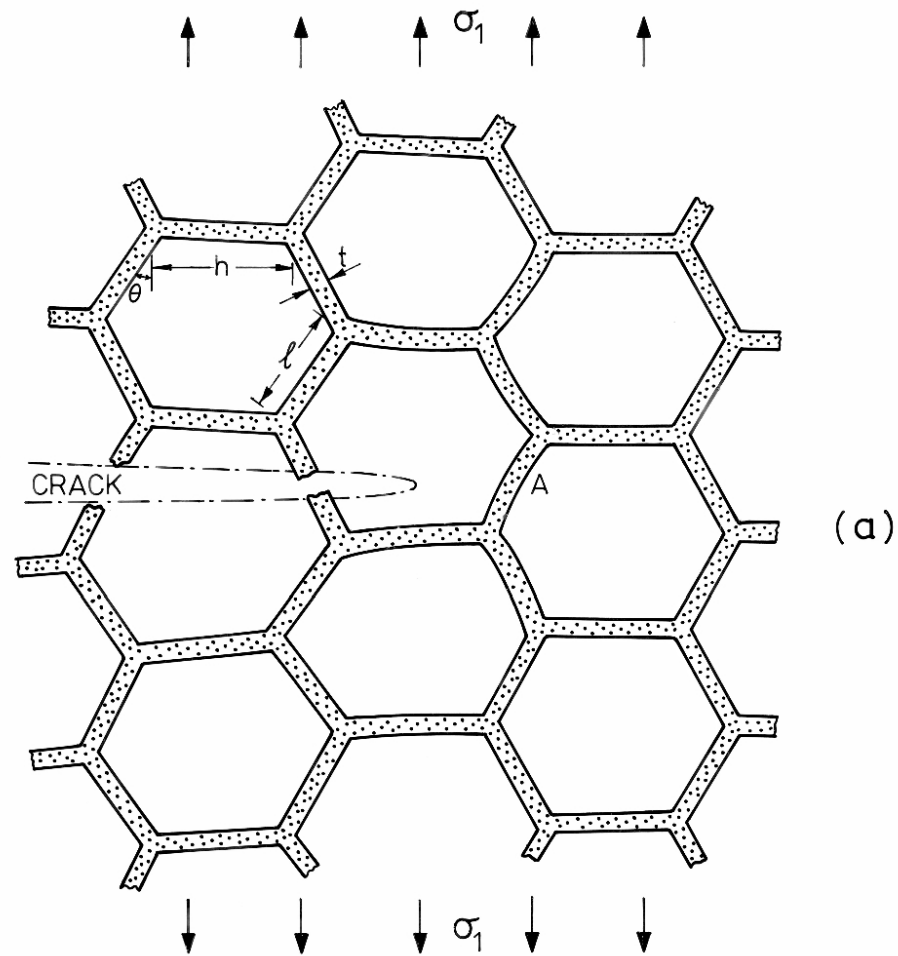
- assume:
- crack length large relative to cell size (continuum assumption)
 - axial forces can be neglected
 - cell wall material has constant modulus of rupture, σ_{fs}

continuum: crack of length $2c$ in a linear elastic solid material normal to a remote tensile stress σ_1 creates a local stress field at the crack tip



$$\sigma_{\text{local}} = \sigma_2 = \frac{\sigma_1 \sqrt{\pi c}}{\sqrt{2\pi r}}$$

Fracture Toughness



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

honey comb: cell walls bent - fail when applied moment = fracture moment

$$M_{app} \propto Pl \quad \text{on wall A} \quad M_f \propto \sigma_{fs} bt^2$$

$$M_{app} \propto Pl \propto \sigma_l l^2 b \propto \frac{\sigma_l \sqrt{c} l^2 b}{\sqrt{l}} \propto \sigma_{fs} bt^2$$

$$(\sigma_f^*)_1 \propto \sigma_{fs} \left(\frac{t}{l}\right)^2 \sqrt{\frac{l}{c}}$$

$$K_{Ic}^* = \sigma_f^* \sqrt{\pi c} = C \sigma_{fs} \left(\frac{t}{l}\right)^2 \sqrt{l}$$

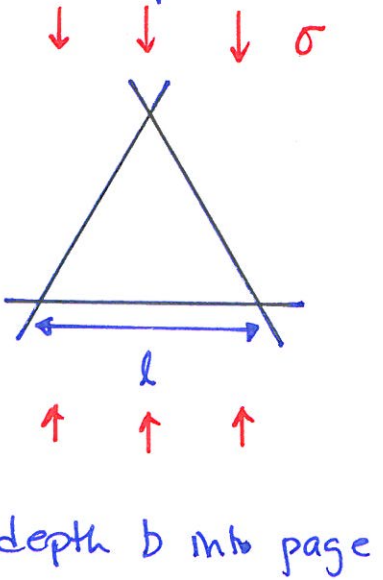
C = constant

depends on cell size, l !

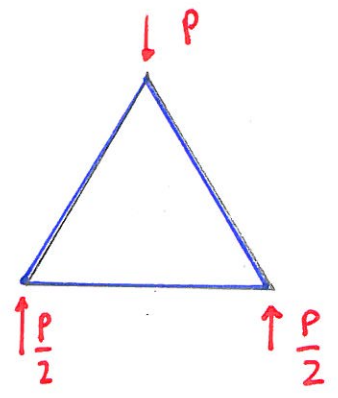
Summary: hexagonal honeycombs, in-plane properties

- linear elastic moduli: E_1^* E_2^* ν_{12}^* G_{12}^*
- plateau stresses $(\sigma_{el}^*)_2$ elastic buckling
(compression)
- σ_{pl}^* plastic yield
- σ_{cr}^* brittle crushing
- fracture toughness K_{Ic}^* brittle fracture
(tension)

Honeycombs: In-plane behaviour - triangular cells



- triangulated structures - trusses
- can analyze as pm-jointed (no moment @ joints)
- forces in members all axial (no bending)
- if joints fixed + include bending, difference ~ 2%
- force in each member proportional to P



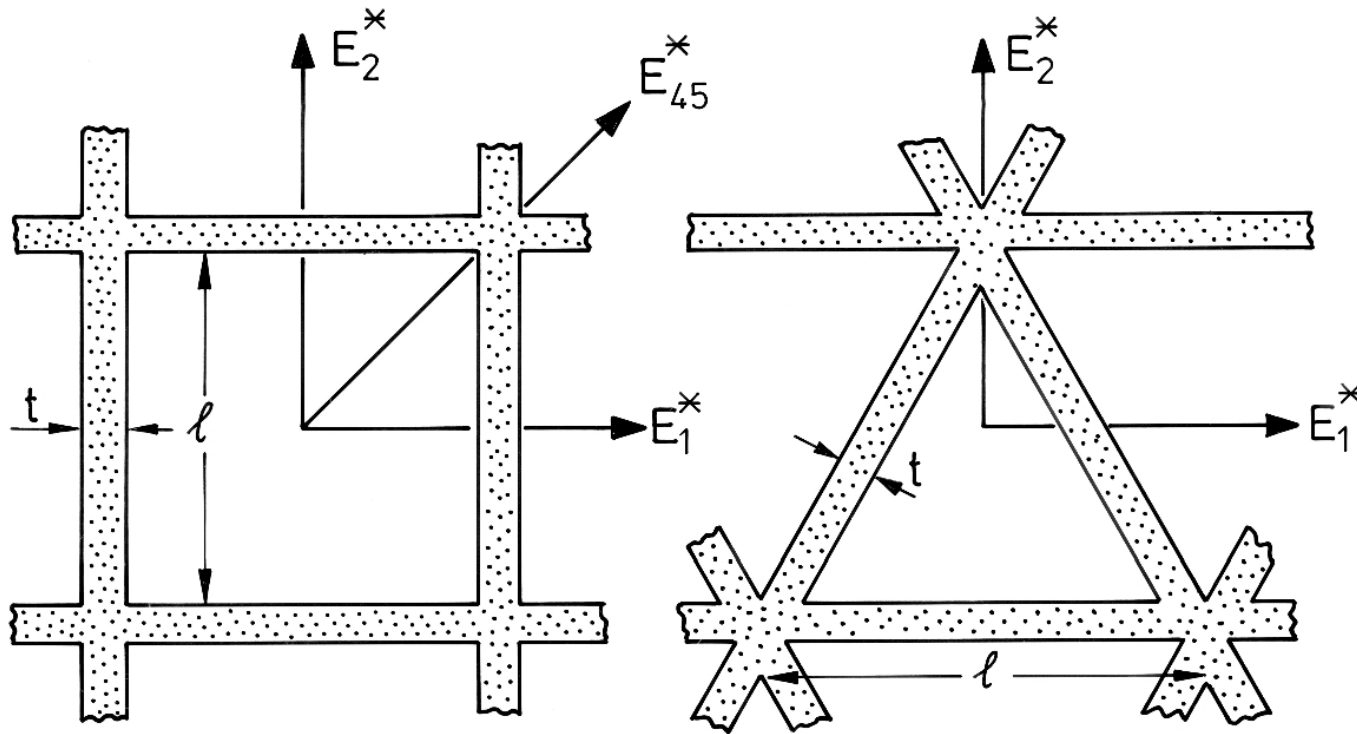
$$\sigma \propto \frac{P}{lb} \quad \epsilon \propto \frac{\delta}{l} \quad \delta \propto \frac{Pl}{AE_s} \quad (\text{axial shortening: Hooke's law})$$

$$E^* \propto \frac{\sigma}{\epsilon} \propto \frac{P}{lb} \frac{l}{\delta} \propto \frac{P}{b} \frac{btE_s}{Pl} \propto E_s \left(\frac{t}{l} \right)$$

$$E^* = c E_s \left(\frac{t}{l} \right)$$

exact calculation: $E^* = 1.15 E_s \left(\frac{t}{l} \right)$ for equilateral triangle

Square and Triangular Honeycombs



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

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