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6.641 Electromagnetic Fields, Forces, and Motion  
Spring 2005

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## Problem 1

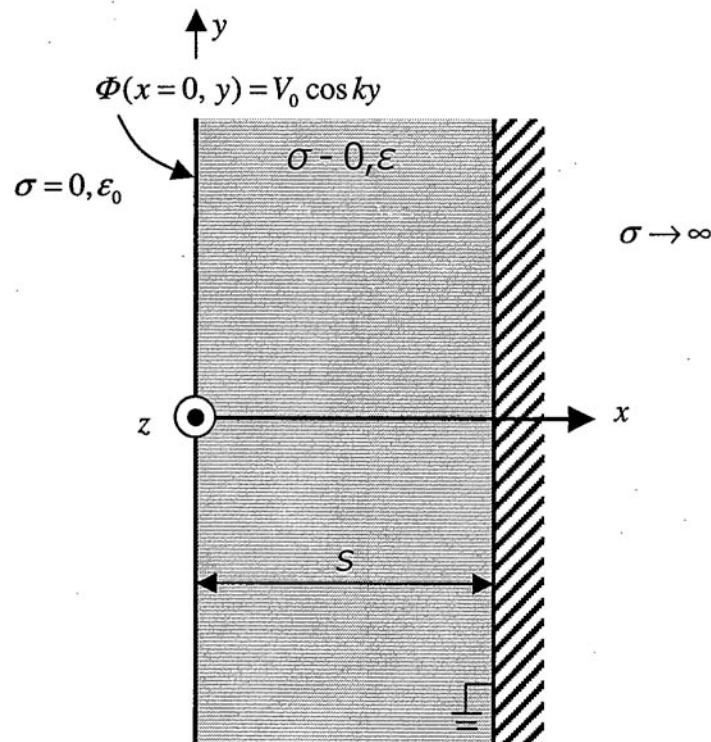


Figure 1: A potential sheet between two lossless dielectrics with a perfect conductor placed at  $x = s$ .

A potential sheet of infinite extent in the  $y$  and  $z$  directions is placed at  $x = 0$  and has potential distribution  $\Phi(x = 0, y) = V_0 \cos(ky)$ . Free space with no conductivity ( $\sigma = 0$ ) and permittivity  $\epsilon_0$  is present for  $x < 0$  while for  $0 < x < s$  a perfectly insulating dielectric ( $\sigma = 0$ ) with permittivity  $\epsilon$  is present. The region for  $x > s$  is a grounded perfect conductor at zero potential.

**A**

**Question:** What are the potential distributions for  $x < 0$  and  $0 < x < s$ ?

**Solution:**

$$\Phi(x, y) = \begin{cases} V_0 \cos(ky) e^{kx} & x < 0 \\ \frac{-V_0 \sinh(k(x-s)) \cos(ky)}{\sinh(ks)} & 0 < x < s \end{cases}$$

**B**

**Question:** What are the surface charge densities at  $x = 0, \sigma_f(x = 0, y)$ , and at  $x = s, \sigma_f(x = s, y)$ ?

**Solution:**

$$E_x = -\frac{\partial\Phi}{\partial x} = \begin{cases} -kV_0 \cos(ky)e^{kx} & x < 0 \\ \frac{kV_0 \cosh(k(x-s)) \cos(ky)}{\sinh(ks)} & 0 < x < s \end{cases}$$

$$\begin{aligned} \sigma_f(x = 0, y) &= \epsilon E_x(x = 0_+, y) - \epsilon_0 E_x(x = 0_-, y) \\ &= [\epsilon_0 + \epsilon \coth(ks)] kV_0 \cos(ky) \\ \sigma_f(x = s, y) &= -\epsilon E_x(x = s_-, y) \\ &= \frac{-\epsilon kV_0 \cos(ky)}{\sinh(ks)} \end{aligned}$$

**C**

**Question:** What is the force, magnitude and direction, on a section of the perfect conductor at  $x = s$  that extends over the region  $0 < y < \frac{\pi}{k}$  and  $0 < z < D$ ?

**Hint:**  $\int \cos^2 y dy = \frac{y}{2} + \frac{\sin(2y)}{4}$ .

**Solution:**

$$\begin{aligned} \frac{F_x}{\text{area}} &= \frac{1}{2} \sigma_f E_x|_{x=s} = -\frac{1}{2} \epsilon E_x^2|_{x=s} = -\frac{1}{2} \frac{\epsilon (kV_0 \cos(ky))^2}{\sinh^2(ks)} \\ F_x &= -\frac{1}{2} \frac{\epsilon k^2 V_0^2 D}{\sinh^2(ks)} \int_0^{\frac{\pi}{k}} \cos^2 ky dy = -\frac{\pi}{4} \frac{\epsilon k V_0^2 D}{\sinh^2(ks)} \end{aligned}$$

## Problem 2

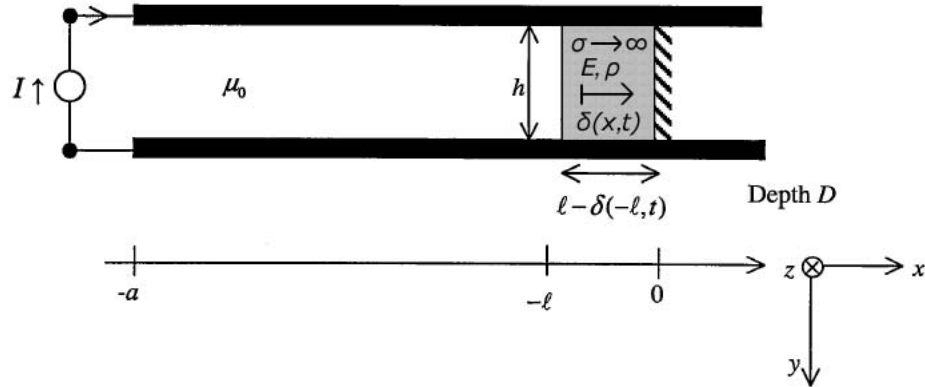


Figure 2: Parallel plate electrodes

Parallel plate electrodes with spacing  $h$  and depth  $D$  are excited by a DC current source  $I$ . An elastic rod surrounded by free space has mass density  $\rho$ , modulus of elasticity  $E$ , equilibrium length  $l$  and has infinite ohmic conductivity  $\sigma$ . The elastic rod end at  $x = 0$  is fixed while the deflections of the rod are described as  $\delta(x, t)$  and are assumed small  $|\delta(x, t)| \ll l$ . The rod width  $l - \delta(-l, t)$  changes as  $I$  is changed because of the magnetic force. The DC current flows as a surface current on the  $x = -(l - \delta(-l, t))$  end of the perfectly conducting rod.

**A**

**Question:** Calculate  $H_z$  in the free space region  $-a < x < -(l - \delta(-l, t))$ . Neglect fringing field effects and assume  $h \ll a$  and  $h \ll D$ .

**Solution:**

$$H_z = \frac{I}{D}$$

**B**

**Question:** Using the Maxwell Stress Tensor calculate the magnetic force per unit area on the  $x = -(l - \delta(-l, t))$  end of the rod.

**Solution:**

$$T_{xx} = \frac{1}{2} \mu_0 [H_x^2 - H_y^2 - H_z^2] \Big|_{x=-(l-\delta(-l,t))} = -\frac{\mu_0}{2} \frac{I^2}{D^2}$$

$$\frac{F_x}{\text{area}} = -T_{xx} \Big|_{x=-(l-\delta(-l,t))} = \frac{\mu_0}{2} \frac{I^2}{D^2}$$

C

**Question:** Calculate the steady state change in rod length  $\delta(x = -l)$ .

**Solution:**

$$\begin{aligned} \rho \frac{\partial^2 \delta}{\partial t^2} &= E \frac{\partial^2 \delta}{\partial x^2} \Rightarrow \delta = ax + b & a &= \frac{-\mu_0 I^2}{2D^2 E} \Rightarrow \delta(x) = \frac{-\mu_0 I^2 x}{2ED^2} \\ \delta(x=0) &= b = 0 \\ E \frac{\partial \delta}{\partial x} \Big|_{x=-l} &= T_{xx} = -\frac{\mu_0 I^2}{2} \frac{1}{D^2} = E_a & \delta(-l) &= \frac{\mu_0 I^2}{2ED^2} \end{aligned}$$

D

**Question:** Noise creates fluctuations  $\delta'(x, t)$  in longitudinal displacement. What are the natural frequencies of the rod?

**Solution:**

$$\begin{aligned} \rho \frac{\partial^2 \delta'}{\partial t^2} &= E \frac{\partial^2 \delta'}{\partial x^2}, \delta'(x, t) = \text{Re} [\hat{\delta}(x) e^{j\omega t}] \\ -\frac{\rho \omega^2}{E} \hat{\delta}(x) &= E \frac{d^2 \hat{\delta}}{dx^2} \Rightarrow \frac{d^2 \hat{\delta}}{dx^2} + k^2 \hat{\delta} = 0, k^2 = \frac{\omega^2 \rho}{E} \\ \hat{\delta}(x) &= A \sin(kx) + B \cos(kx) \\ \hat{\delta}(x=0) &= B = 0 \\ \frac{d\hat{\delta}}{dx} \Big|_{x=-l} &= 0 = kA \cos(kl) \Rightarrow kl = (2n+1) \frac{\pi}{2}, n = 0, 1, 2 \\ \omega_n \sqrt{\frac{\rho}{E}} &= k_n = (2n+1) \frac{\pi}{2l} \\ \omega_n &= \sqrt{\frac{E}{\rho}} \frac{\pi}{2l} (2n+1), n = 0, 1, 2 \end{aligned}$$

### Problem 3

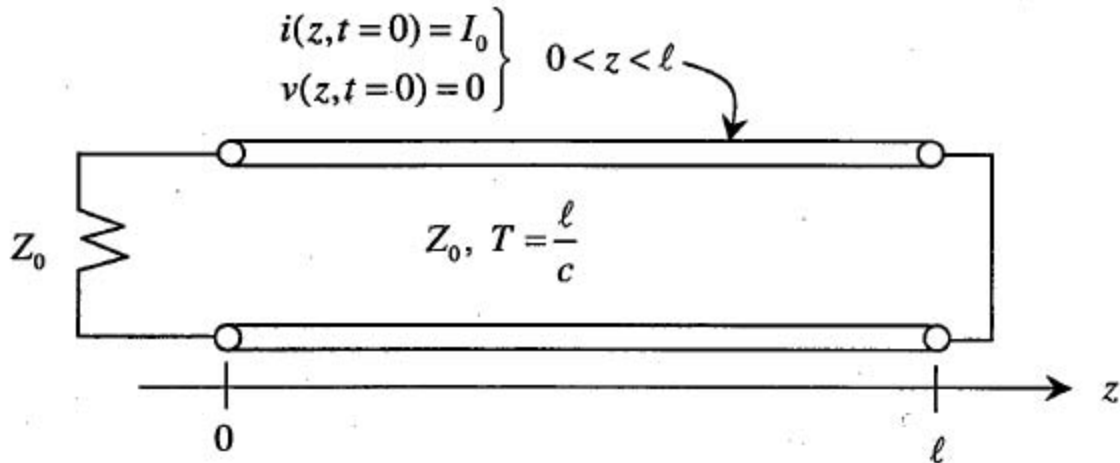


Figure 3: An electrical transmission line

An electrical transmission line of length  $l$  has characteristic impedance  $Z_0$ . Electromagnetic waves can travel on the line at speed  $c$ , so that the time to travel one-way over the line length  $l$  is  $T = \frac{l}{c}$ . The line is matched at  $z = 0$  and is short circuited at  $z = l$ . At time  $t = 0$ , a lightning bolt strikes the entire line so that there is a uniform current along the line but with zero voltage:

$$\begin{aligned} i(z, t = 0) &= I_0 & 0 < z < l \\ v(z, t = 0) &= 0 & 0 < z < l \end{aligned}$$

Since the voltage and current obey the telegrapher's relations:

$$\begin{aligned} \frac{\partial v}{\partial z} &= -L \frac{\partial i}{\partial t}, c = \frac{1}{\sqrt{LC}} \\ \frac{\partial i}{\partial z} &= -C \frac{\partial v}{\partial t}, Z_0 = \sqrt{\frac{L}{C}} \end{aligned}$$

the voltage and current along the line are related as

$$\begin{aligned} v + iZ_0 &= c_+ \text{ on } \frac{dz}{dt} = c \Rightarrow v = \frac{c_+ + c_-}{2} \\ v - iZ_0 &= c_- \text{ on } \frac{dz}{dt} = -c \Rightarrow iZ_0 = \frac{c_+ - c_-}{2} \end{aligned}$$

**A**

**Question:** The solutions for  $v(z, t)$  and  $i(z, t)$  can be found using the method of characteristics within each region shown below (see exam questions). Within regions 1-9 give the values of  $c_+, c_-, v$  and  $iZ_0$ .

**Solution:** See Figure 4.

**B**

**Question:** Plot  $v(z, t = \frac{T}{4})$  and  $i(z, t = \frac{T}{4})$ .

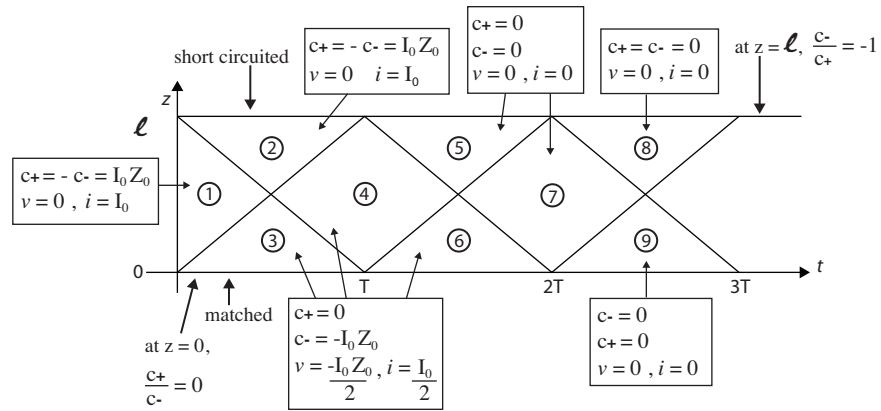


Figure 4: Values of  $c_+$ ,  $c_-$ ,  $v$  and  $iZ_0$  in regions 1-9. (Image by MIT OpenCourseWare.)

**Solution:**

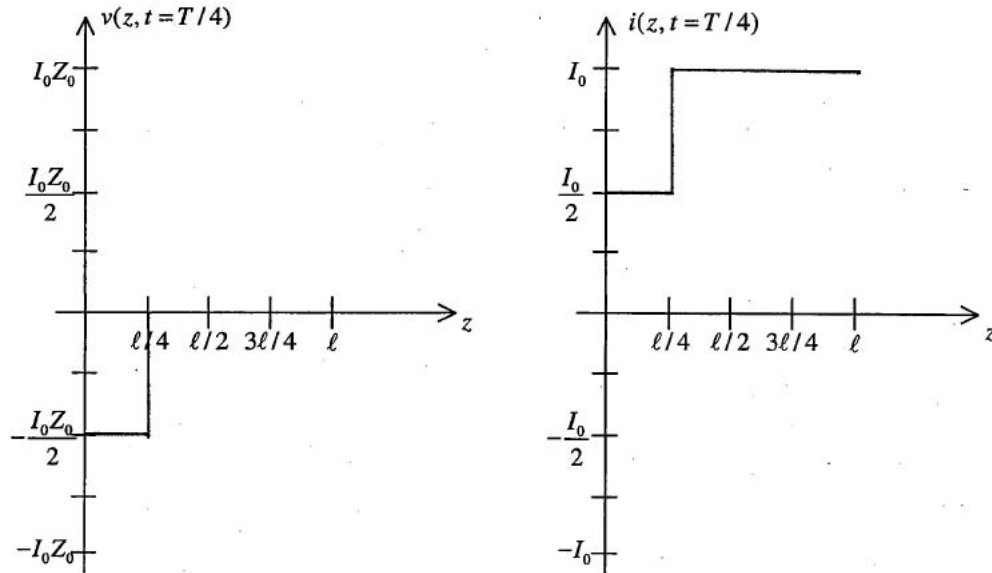


Figure 5: plots of  $v(z, t = \frac{T}{4})$  and  $i(z, t = \frac{T}{4})$  for problem 2, part B (Image by MIT OpenCourseWare.)

**C**

**Question:** How long a time does it take for the transmission line to have  $v(z, t) = 0$  and  $i(z, t) = 0$  everywhere for  $0 < z < l$  for all further time?

**Solution:**  $2T$

## Problem 4

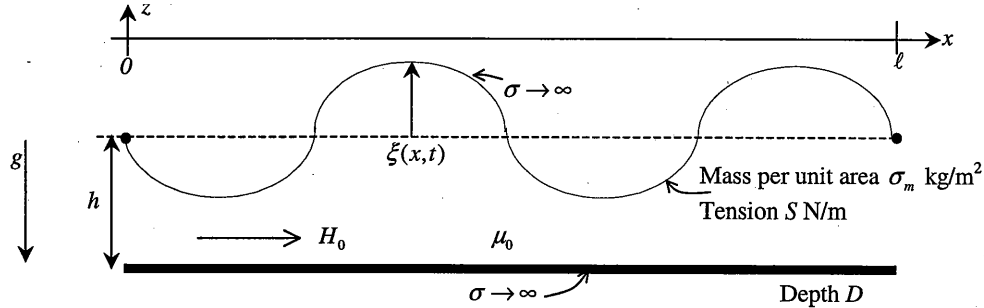


Figure 6: A perfectly conducting membrane stressed from below by magnetic field  $H_0 \bar{i}_x$

A perfectly conducting membrane of depth  $D$  with mass per unit area  $\sigma_m$  and tension  $S$  is a distance  $h$  above a rigid perfect conductor. The membrane and rigid conductor are in free space and support currents such that when the membrane is flat,  $\xi(x, t) = 0$ , the static uniform magnetic field intensity is  $H_0$ . As the membrane deforms, the flux through the region between membrane and rigid conductor is conserved. The system is in a downward gravity field with gravitational acceleration  $\bar{g} = -g\bar{i}_z$ . The membrane deflection has no dependence on  $y$  and is fixed at its two ends at  $x = 0$  and  $x = l$ .

### A

**Question:** Assuming that  $\xi(x, t) \ll h$  and that the only significant magnetic field component is  $x$  directed, how is  $H_x(x, t)$  approximately related to  $\xi(x, t)$  to linear terms in  $\xi(x, t)$ ?

**Solution:**

$$H_x(h + \xi) = H_0 h \Rightarrow H_x = \frac{H_0 h}{h + \xi} = \frac{H_0}{1 + \frac{\xi}{h}} \approx H_0 \left(1 - \frac{\xi}{h}\right)$$

### B

**Question:** Using the Maxwell Stress tensor and the result of part(a), to linear terms in small displacement  $\xi(x, t)$ , what is the  $z$  directed magnetic force per unit area,  $F_z$ , on the membrane?

**Solution:**

$$T_{zz} = \frac{\mu_0}{2} \left( H_z^2 - H_x^2 - H_y^2 \right) = -\frac{\mu_0}{2} H_0^2 \left(1 - \frac{\xi}{h}\right)^2 \approx -\frac{\mu_0 H_0^2}{2} \left(1 - \frac{2\xi}{h}\right)$$

$$F_z = -T_{zz} = \frac{\mu_0 H_0^2}{2} \left(1 - \frac{2\xi}{h}\right)$$



**C**

**Question:** To linear terms in small displacement  $\xi(x, t)$ , express the membrane equation of motion in the form

$$a \frac{\partial^2 \xi}{\partial t^2} = b \frac{\partial^2 \xi}{\partial x^2} + c \xi + d$$

What are  $a, b, c$ , and  $d$ ?

**Solution:**

$$\begin{aligned} \sigma_m \frac{\partial^2 \xi}{\partial t^2} &= S \frac{\partial^2 \xi}{\partial x^2} + F_z - \sigma_m g \\ &= S \frac{\partial^2 \xi}{\partial x^2} + \mu_0 \frac{H_0^2}{2} \left( 1 - \frac{2\xi}{h} \right) - \sigma_m g \end{aligned}$$

$$a = \sigma_m, b = S, c = -\frac{\mu_0 H_0^2}{h}, d = \frac{\mu_0 H_0^2}{2} - \sigma_m g$$

**D**

**Question:** What value of  $H_0$  is needed so that in static equilibrium the membrane has no sag,  $\xi(x, t) = 0$ .

**Solution:**

$$\frac{\mu_0 H_0^2}{2} = \sigma_m g \Rightarrow H_0 = \left[ \frac{2\sigma_m g}{\mu_0} \right]^{\frac{1}{2}}$$

**E**

**Question:** About the equilibrium of part (d), what is the  $\omega-k$  dispersion relation for membrane deflections of the form

$$\xi(x, t) = \text{Re} \left[ \hat{\xi} e^{j(\omega t - kx)} \right]?$$

Solve for  $k$  as a function of  $\omega$  and system parameters.

**Solution:**

$$\begin{aligned} -\sigma_m \omega^2 &= -S k^2 - \frac{\mu_0 H_0^2}{h} \\ k^2 &= \frac{\sigma_m \omega^2}{S} - \frac{\mu_0 H_0^2}{hS} \\ k &= \pm \left[ \frac{\sigma_m \omega^2}{S} - \frac{\mu_0 H_0^2}{hS} \right]^{\frac{1}{2}} = \pm k_0 \\ k_0 &= + \left[ \frac{\sigma_m \omega^2}{S} - \frac{\mu_0 H_0^2}{hS} \right]^{\frac{1}{2}} \end{aligned}$$

**F**

**Question:** Using all the values of  $k$  found in part (e), find a superposition of solutions of the form of  $\xi(x, t)$  given in (e) that satisfy the zero deflection boundary conditions at the ends of the membrane at  $x = 0$  and  $x = l$ . What are the allowed values of  $k$ ?

**Solution:**

$$\begin{aligned}\xi(x, t) &= \operatorname{Re} \left[ e^{j\omega t} \left[ \hat{\xi}_1 e^{-jk_0 x} + \hat{\xi}_2 e^{+jk_0 x} \right] \right] \\ \xi(x = 0, t) = 0 &= \operatorname{Re} \left[ e^{j\omega t} \left[ \hat{\xi}_1 + \hat{\xi}_2 \right] \right] \Rightarrow \hat{\xi}_2 = -\hat{\xi}_1 \\ \xi(x, t) &= \operatorname{Re} \left[ e^{j\omega t} \hat{\xi}_1 \left[ e^{-jk_0 x} - e^{jk_0 x} \right] \right] \\ &= \operatorname{Re} \left[ e^{j\omega t} \hat{\xi}_1 (-2j) \sin(k_0 x) \right] \\ \xi(x = l, t) = 0 &= \operatorname{Re} \left[ e^{j\omega t} \hat{\xi}_1 (-2j) \sin(k_0 l) \right] = 0 \\ \sin(k_0 l) = 0 &\Rightarrow k_0 l = n\pi, n = 1, 2, \dots\end{aligned}$$

**G**

**Question:** Is this system always stable or under what conditions can it be unstable? When stable, what are the natural frequencies and if unstable what are the growth rates of the instability?

**Solution:**

$$\begin{aligned}\omega^2 &= \frac{S}{\sigma_m} k^2 + \frac{\mu_0 H_0^2}{h\sigma_m} \\ \omega_n &= \left[ \frac{S}{\sigma_m} \left( \frac{n\pi}{l} \right)^2 + \frac{\mu_0 H_0^2}{h\sigma_m} \right]^{\frac{1}{2}} \text{ Always stable as } \omega_n \text{ real.}\end{aligned}$$