

10.675 LECTURE 6

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1. TODAY

- Variational Principle
- Derive HF (Hartree-Fock) Equations
- Interpretation of Solutions to HF equations

2. VARIATIONAL PRINCIPLE

Idea: The closer our guess of C_i 's of Ψ_{trial} , the lower our energy.

$\tilde{\Phi} \Rightarrow$ trial function

$$E[\tilde{\Phi}] = \langle \tilde{\Phi} | H | \tilde{\Phi} \rangle$$

Functional Definition : maps functions to numbers

3. FUNCTIONAL VARIATION

vary $\tilde{\Phi} \rightarrow \tilde{\Phi} + \delta\tilde{\Phi}$

substitute in and solve

$$E[\tilde{\Phi} + \delta\tilde{\Phi}] = \langle \tilde{\Phi} + \delta\tilde{\Phi} | H | \tilde{\Phi} + \delta\tilde{\Phi} \rangle$$

expand this via linear first order term.

$$E[\tilde{\Phi} + \delta\tilde{\Phi}] = E[\tilde{\Phi}] + \langle \delta\tilde{\Phi} | H | \tilde{\Phi} \rangle + \langle \tilde{\Phi} | H | \delta\tilde{\Phi} \rangle + \text{higher order terms.}$$

$$E[\tilde{\Phi} + \delta\tilde{\Phi}] = E[\tilde{\Phi}] + \delta E$$

Trying to approach the true solution such that $H|\Phi\rangle = E|\Phi\rangle \Rightarrow \delta E = 0$

Normally, E will always be a min, so we don't have to worry about local/global max solutions.

There exists an infinite number of solutions.

$$H|\tilde{\Phi}_\alpha\rangle = E_\alpha|\tilde{\Phi}_\alpha\rangle$$

$$\alpha = 0, 1, 2, \dots$$

$$E_0 \leq E_1 \leq E_2 \dots$$

H is hermitian

$$H = H^\dagger, \text{ and } E_\alpha \text{'s are real}$$

$\tilde{\Phi}_\alpha$'s are orthonormal.

$$\langle \tilde{\Phi}_\alpha | \tilde{\Phi}_\beta \rangle = \delta_{\alpha\beta}$$

$$\langle \tilde{\Phi}_\alpha | H | \tilde{\Phi}_\beta \rangle = E_\alpha \delta_{\alpha\beta}$$

4. EXPANSION

Expand trial wave function $\tilde{\Phi}$ in terms of the eigenstates Φ .

$$|\tilde{\Phi}\rangle = \sum_\alpha C_\alpha |\Phi_\alpha\rangle$$

$$|\tilde{\Phi}\rangle = \sum_\alpha |\Phi_\alpha\rangle \langle \Phi_\alpha | \tilde{\Phi} \rangle$$

$$\langle \tilde{\Phi} | = \sum_\alpha C_\alpha^* \langle \Phi_\alpha |$$

Date: Fall 2004.

$$\begin{aligned}
\langle \tilde{\Phi} | &= \sum_{\alpha} \langle \tilde{\Phi} | \Phi_{\alpha} \rangle \langle \Phi_{\alpha} | \\
\langle \tilde{\Phi} | \tilde{\Phi} \rangle &= 1 = \sum_{\alpha} \langle \tilde{\Phi} | \Phi_{\alpha} \rangle \langle \Phi_{\alpha} | \tilde{\Phi} \rangle = \sum_{\alpha} | \langle \Phi_{\alpha} | \tilde{\Phi} \rangle |^2 \\
\text{and} \\
\langle \tilde{\Phi} | H | \tilde{\Phi} \rangle &= \sum_{\alpha\beta} \langle \tilde{\Phi} | \Phi_{\alpha} \rangle \langle \Phi_{\alpha} | H | \Phi_{\beta} \rangle \langle \Phi_{\beta} | \tilde{\Phi} \rangle = \sum_{\alpha} E_{\alpha} | \langle \Phi_{\alpha} | \tilde{\Phi} \rangle |^2 \\
E[\Phi] &= \frac{\langle \tilde{\Phi} | H | \tilde{\Phi} \rangle}{\langle \tilde{\Phi} | \tilde{\Phi} \rangle} = \frac{\sum_{\alpha} E_{\alpha} | \langle \Phi_{\alpha} | \tilde{\Phi} \rangle |^2}{\sum_{\alpha} | \langle \Phi_{\alpha} | \tilde{\Phi} \rangle |^2} = \frac{\sum_{\alpha} E_{\alpha} | C_{\alpha} |^2}{\sum_{\alpha} | C_{\alpha} |^2} \geq E_o \\
\text{Use variational principle to derive HF equations given the determinant} \\
|\Psi_o\rangle &= |\chi_1, \chi_2, \chi_3, \dots, \chi_N\rangle \\
E_o &= \langle \Psi_o | H | \Psi_o \rangle = E_o[\{\chi_{\alpha}\}] \\
\text{Minimize with respect to } \chi_{\alpha} \text{'s functions w/constraint of } \chi_{\alpha} \text{'s being orthonormal.} \\
\int d\vec{x} \chi_a^*(1) \chi_b(1) &= [a|b] = \delta_{ab}
\end{aligned}$$

5. LAGRANGE METHOD OF UNDETERMINED MULTIPLIERS

$$\begin{aligned}
L[\{\chi_{\alpha}\}] &= E_o[\{\chi_{\alpha}\}] - \sum_{a=1}^N \sum_{b=1}^N \epsilon_{ab} ([a|b] - \delta_{ab}) \\
\text{Where } \epsilon_{ab} &\text{ are the undetermined multipliers.} \\
\frac{\delta L}{\delta \epsilon_{ab}} &= 0 \\
\frac{\delta L}{\delta \chi_a} &= 0 \\
\text{vary: } \chi_a &\rightarrow \chi_a + \delta \chi_a \\
\delta L &= 0 \delta E_o - \sum_a^N \sum_b^N \epsilon_{ab} \delta [a|b] = 0 \\
\delta [a|b] &= [\delta \chi_a | \chi_b] + [\chi_a | \delta \chi_b] \\
\delta E_o &= \sum_{a=q}^N ([\delta \chi_a | h | \chi_a] + [\chi_a | h | \delta \chi_a] \\
&+ \frac{1}{2} \sum_a \sum_b ([\delta \chi_a \chi_a | \chi_b \chi_b] + [\chi_a \delta \chi_a | \chi_b \chi_b] + [\chi_a \chi_a | \delta \chi_b \chi_b] + [\chi_a \chi_a | \chi_b \delta \chi_b]) \\
&+ \frac{1}{2} \sum_a \sum_b ([\delta \chi_a \chi_b | \chi_b \chi_a] + [\chi_a \delta \chi_b | \chi_b \chi_a] + [\chi_a \chi_b | \delta \chi_b \chi_a] + [\chi_a \chi_b | \chi_b \delta \chi_a]) \\
&= \sum_a^N [\delta \chi_a | h | \chi_a] + \sum_a \sum_b ([\delta \chi_a \chi_b | \chi_b \chi_b] - [\delta \chi_a \chi_b | \chi_b \chi_a]) + \text{the complex} \\
&\text{conjugate terms.} \\
\delta L = 0 &= \sum_1^N [\delta \chi_a | 1 | \chi_a] + \sum_{ab} [\delta \chi_a \chi_a | \chi_b \chi_b] - [\delta \chi_a \chi_b | \chi_b \chi_a] - \sum [ab \epsilon_a [\delta \chi_a = \chi_b] + \\
&\text{the complex conjugate terms.} \\
\delta L &= \sum_{a=1}^N \int dx \delta \chi_a(1) [h(1) \chi_a(1) + \sum_{b=1}^N (J_b(1) - K_b(1) \chi_a(1) - \sum_{b=1}^N \epsilon_{ba} \chi_b(1))] = 0 \\
&+ \text{Complex conjugate terms.}
\end{aligned}$$

6. CONTINUED...

$$\begin{aligned}
[h(1) + \sum_b^N J_b(1) - K_b(1)] \chi_a(1) &= \sum_b^N \epsilon_{ab} \chi_b(1) \\
\text{"Canonical Form" via unitary transformation to diagonalize } f|\chi_a\rangle &= \epsilon|\chi_a\rangle \\
\text{Canonical HF equations.} \\
\epsilon_a &= \langle \chi_a | f | \chi_a \rangle = \langle \chi_a | h + \sum_b (J_b - K_b) | \chi_a \rangle \\
&= \langle a | h | a \rangle + \sum_b \langle ab | ab \rangle - \langle ab | ba \rangle \\
&= \langle a | h | a \rangle + \sum_b \langle ab | ab \rangle \text{ where } |a\rangle \text{'s are occupied spin orbitals} \\
\epsilon_r &= \langle r | h | r \rangle + \sum_b \langle rb | rb \rangle \text{ where } |r\rangle \text{'s are "virtual" spin orbitals} \\
\sum_a^N \epsilon_a &= \sum_a^N \langle a | h | a \rangle + \sum_a^N \sum_b^N \langle ab | ab \rangle \\
E_o \sum_a^N \langle a | h | a \rangle &+ 1/2 \sum \sum \langle ab | ab \rangle \neq \sum_a^N \epsilon_a \\
\text{because of the coefficients on the double summation.} \\
\text{Example: } \epsilon_a &\text{ includes interactions with all other orbitals, as does } \epsilon_b, \text{ so they are} \\
&\text{being double counted in SCF.}
\end{aligned}$$

7. SIGNIFICANCE OF ORBITAL ENERGIES (Koopman's THEOREM)

$$\begin{aligned}
E_o^N &= \sum_a \langle a | h | a \rangle + 1/2 \sum_a \sum_b \langle ab | ab \rangle \\
\text{when you take out one electron}
\end{aligned}$$

$$E_c^{N-1} = \sum_{anec} \langle a|h|a \rangle + 1/2 \sum_{anec} \sum_{bnec} \langle ab||ab \rangle$$

$$\text{Ionization potential} = E_c^{N-1} - E_o^N$$

$$= - \langle c|h|c \rangle - 1/2 \sum_a \langle ac||ac \rangle - 1/2 \sum_b \langle cb||cb \rangle$$

$$= - \langle c|h|c \rangle - \sum_b \langle cb||cb \rangle = \epsilon_c$$

"Electron Affinity" = $E_A = E_o^N - E_r^{N+1} = \epsilon_r$ = the energy of the system with an additional electron in the virtual levels.

This is the "frozen" orbital approximation.

$$\begin{array}{ccc}
 & & \uparrow\downarrow \\
 & & \leftarrow \epsilon_r \\
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 \uparrow\downarrow & \uparrow\downarrow & \\
 \uparrow\downarrow & \uparrow\downarrow & \\
 \uparrow\downarrow & \uparrow\downarrow & \\
 \uparrow\downarrow & - & \leftarrow \epsilon_a \\
 \uparrow\downarrow & \uparrow\downarrow &
 \end{array}$$